

## Micro-Macro Modelling and Simulation of Liquid-Vapour Flow



# A Multiscale Method for Compressible Liquid-Vapor Flow with Surface Tension

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# Free Boundary Value Problem in $\mathbb{R}^d$

## Bulk Phases

$$\begin{aligned}\varrho_t + \operatorname{div}(\varrho \mathbf{v}) &= 0, \\ (\varrho \mathbf{v})_t + \operatorname{div}(\varrho \mathbf{v} \otimes \mathbf{v} + p(\varrho) \mathbf{I}) &= \mathbf{0}.\end{aligned}$$

## Phase Boundary

$$\begin{aligned}\llbracket \varrho(\mathbf{v} \cdot \mathbf{n} - \sigma) \rrbracket &= 0, \\ \llbracket \varrho(\mathbf{v} \cdot \mathbf{n} - \sigma) \mathbf{v} + p(\varrho) \mathbf{n} \rrbracket &= (d-1)\gamma\kappa \mathbf{n}.\end{aligned}$$

+ condition for entropy dissipation

For a well-posedness result with zero  
entropy dissipation see  
[Benzoni-Gavage, Freistühler 2004]

# Outline

Heterogeneous Multiscale Method

Microscale

Sharp Interface Model

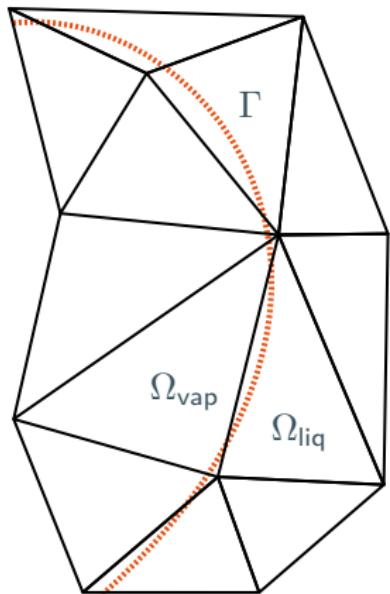
A Curvature Dependent Pressure Function

Solution

Application of the HMM in 2d

Outlook

# Heterogeneous Multiscale Method



Domain  $\Omega \in \mathbb{R}^d$  with fluid in two phases at time  $t > 0$

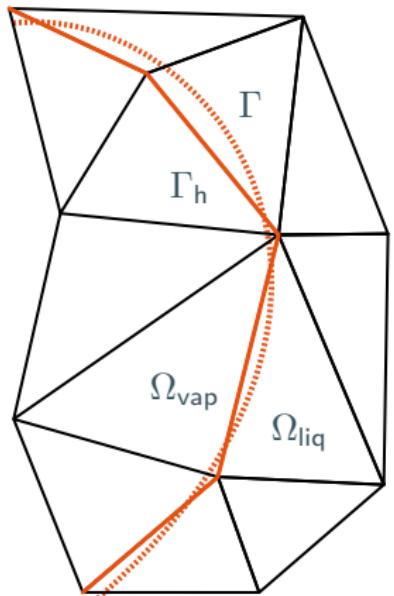
- liquid  $\Omega_{\text{liq}}(t)$ ,
- vapour  $\Omega_{\text{vap}}(t)$ ,
- curved phase boundary  $\Gamma(t) = \partial\Omega_{\text{liq}}(t) \cap \partial\Omega_{\text{vap}}(t)$ .

We consider physical quantities in Euler coordinates  $(x, t)$

- density  $\varrho(x, t) > 0$ ,
- velocity  $v(x, t) \in \mathbb{R}^d$ .

[Engquist et al. 2007]

# Heterogeneous Multiscale Method



Domain  $\Omega \in \mathbb{R}^d$  with fluid in two phases at time  $t > 0$

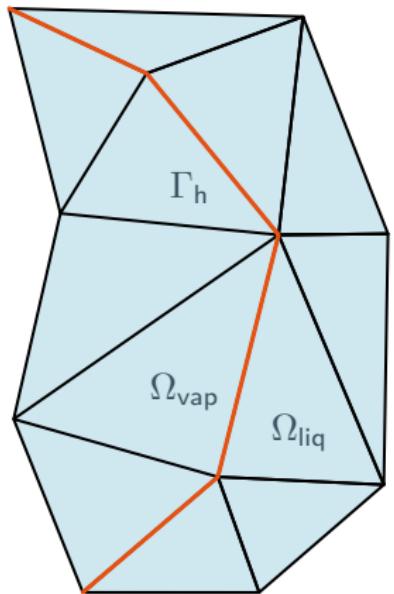
- liquid  $\Omega_{\text{liq}}(t)$ ,
- vapour  $\Omega_{\text{vap}}(t)$ ,
- curved phase boundary  $\Gamma(t) = \partial\Omega_{\text{liq}}(t) \cap \partial\Omega_{\text{vap}}(t)$ .

## Phase Boundary

Approximate  $\Gamma(t)$  with  $\Gamma_h(t)$  consisting of element edges.

[Engquist et al. 2007]

# Heterogeneous Multiscale Method



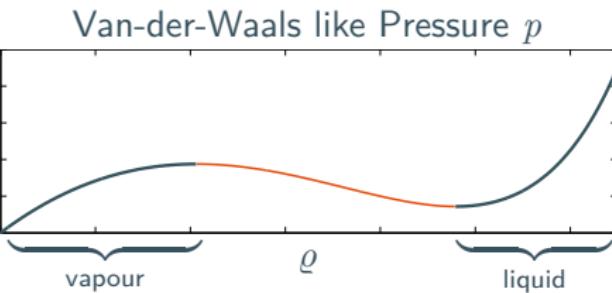
[Engquist et al. 2007]

## Macroscale

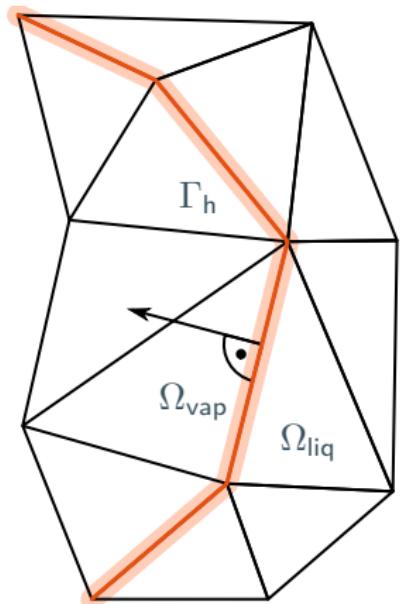
Dynamics of the bulk phases  $\Omega_{\text{liq}}$  and  $\Omega_{\text{vap}}$ .  
Model: Isothermal Euler equation

$$\begin{aligned}\varrho_t + \operatorname{div}(\varrho \mathbf{v}) &= 0, \\ (\varrho \mathbf{v})_t + \operatorname{div}(\varrho \mathbf{v} \otimes \mathbf{v} + p(\varrho) \mathbf{I}) &= \mathbf{0},\end{aligned}$$

with a pressure function  $p$  that covers liquid and vapour phase.



# Heterogeneous Multiscale Method



## Microscale

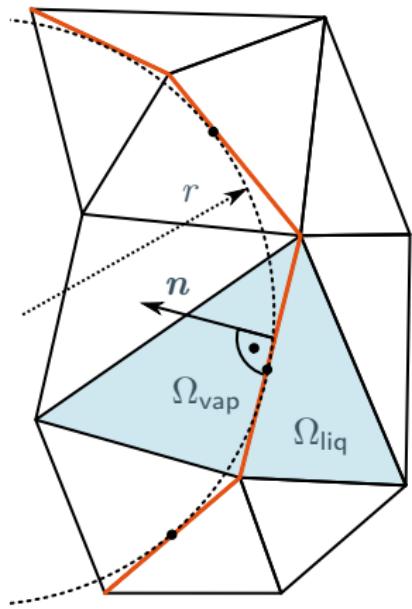
Dynamics at the phase boundary.

We assume that it is sufficient to consider a 1D Riemann type problem on the microscale.

- sharp interface
  - isothermal Euler equation and jump conditions across the phase boundary
- diffuse interface
  - Navier-Stokes-Korteweg model

[Engquist et al. 2007]

# Heterogeneous Multiscale Method



## Transfer Operator

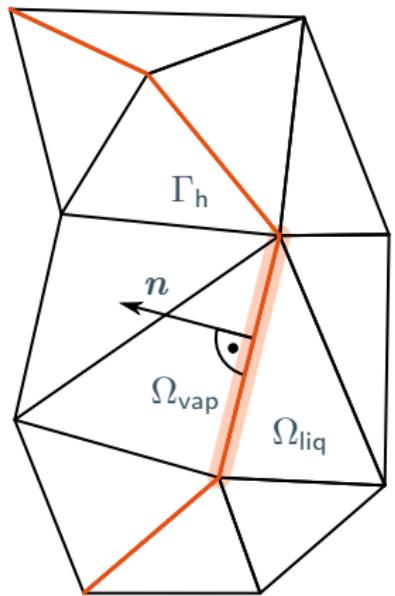
Macro- to microscale

- data reconstruction
  - initial states for microscale problems
- curvature reconstruction
  - applying  $\Gamma_h$  (figure:  $\kappa = \frac{1}{r}$ )
  - level set function

$$\kappa = \operatorname{div} \left( \frac{\operatorname{grad} \varphi}{|\operatorname{grad} \varphi|} \right)$$

[Engquist et al. 2007]

# Heterogeneous Multiscale Method



## Transfer Operator

Micro- to macroscale

- data compression
  - fluxes for phase boundary edges
- interface tracking
  - level set function

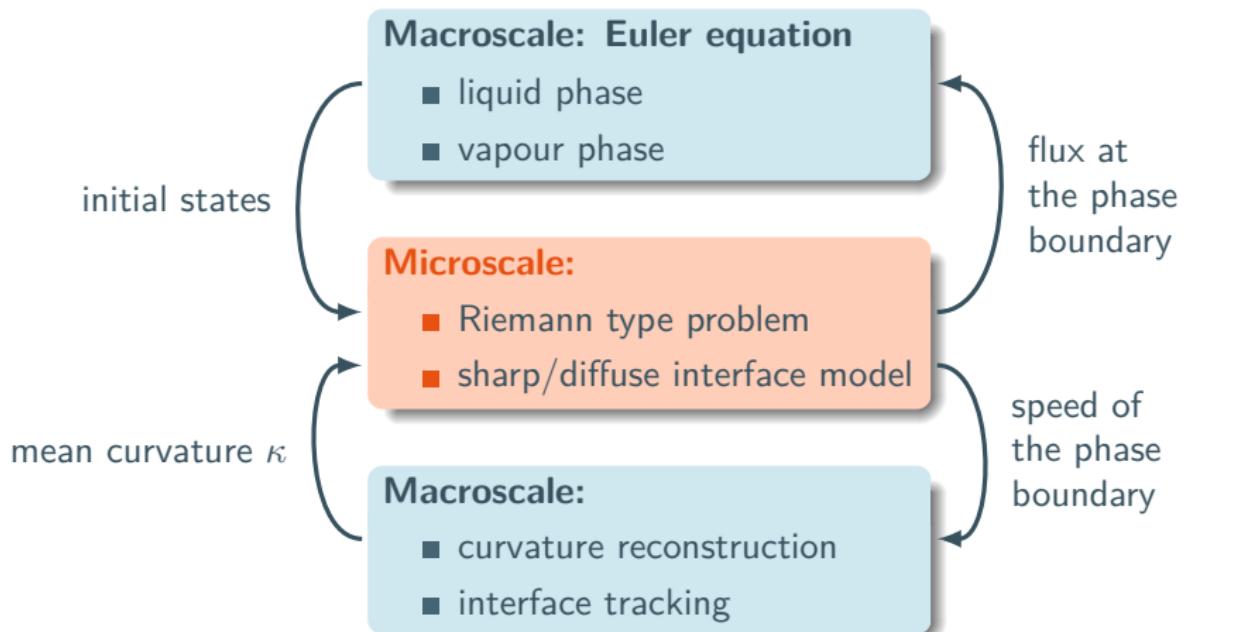
$$\frac{\partial \varphi}{\partial t} + (\sigma \mathbf{n}) \cdot \operatorname{grad} \varphi = \mathbf{0}$$

[Engquist et al. 2007]

# Heterogeneous Multiscale Method

## Algorithm

### Reconstruction:



## Microscale

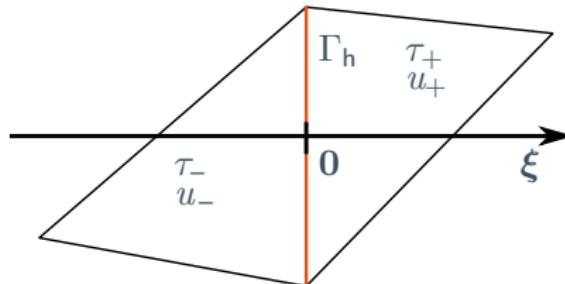
## Sharp Interface Model

Within a macro time step, there are essentially Riemann problems to solve:

$$(\tau, u)(\xi, 0) := \begin{cases} (\tau_-, u_-) & \text{for } \xi \leq 0, \\ (\tau_+, u_+) & \text{for } \xi > 0. \end{cases} \quad (\text{RP})$$

In Lagrange coordinates  $(\xi, t)$ :

- specific volume  
 $\tau := \varrho^{-1} > 0$
- velocity  $u \in \mathbb{R}$



## Microscale

Sharp Interface Model

$$\begin{pmatrix} \tau \\ u \end{pmatrix}_t + \begin{pmatrix} -u \\ \tilde{p}(\tau) \end{pmatrix}_\xi = \begin{pmatrix} 0 \\ 0 \end{pmatrix} \quad \text{in } \mathbb{R} \times (0, \infty). \quad (\text{Euler})$$

We solve the Riemann problem (RP) for the isothermal Euler equation (Euler) such that

1. the solution in bulk phases is entropy solution,
2. there exists one phase boundary that satisfies

$$\begin{aligned} \sigma [\![\tau]\!] + [\![u]\!] &= 0, \\ \sigma [\![u]\!] + [\![\tilde{p}(\tau)]!] &= (d-1)\gamma\kappa, \end{aligned} \quad (\text{RH}^\kappa)$$

3. Liu's entropy criterion holds ([Liu 1974]).

$\sigma$	speed of the phase boundary	$d$	spatial dimension
$\gamma > 0$	surface tension coefficient	$\kappa$	mean curvature
$\tilde{p}(\tau)$	$:= p(\tau^{-1})$ Van-der-Waals Pressure	$[\![\tau]\!]$	$:= \tau_- - \tau_+$

# Microscale

## A Curvature Dependent Pressure Function

We introduce a curvature dependent pressure function  $\tilde{p}^\kappa$ :

$$\tilde{p}^\kappa(\tau) := \begin{cases} \tilde{p}(\tau) + (d-1)\gamma\kappa & \text{in the liquid phase,} \\ \tilde{p}(\tau) & \text{in the vapour phase.} \end{cases}$$

$$\begin{pmatrix} \tau \\ u \end{pmatrix}_t + \begin{pmatrix} -u \\ \tilde{p}^\kappa(\tau) \end{pmatrix}_\xi = \begin{pmatrix} 0 \\ 0 \end{pmatrix} \quad \text{in } \mathbb{R} \times (0, \infty) \quad (\text{Euler}^\kappa)$$

The Rankine Hugoniot conditions for  $(\text{Euler}^\kappa)$  coincides with  $(\text{RH}^\kappa)$ .

Appropriate free energy

$$\tilde{\psi}^\kappa(\tau) := \begin{cases} \tilde{\psi}(\tau) - (d-1)\tau\gamma\kappa & \text{in the liquid phase,} \\ \tilde{\psi}(\tau) & \text{in the vapour phase.} \end{cases}$$

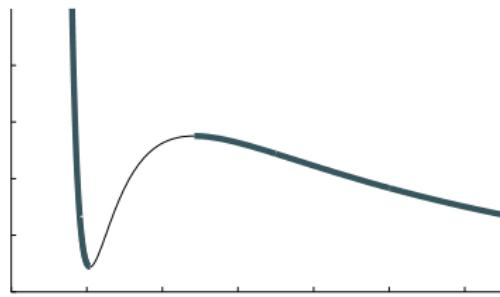
such that  $\frac{d}{d\tau}\tilde{\psi}^\kappa(\tau) = -\tilde{p}^\kappa(\tau)$  and  $\tilde{\psi}^\kappa + \tau \tilde{p}^\kappa = \tilde{\psi} + \tau \tilde{p}$ .

# Microscale

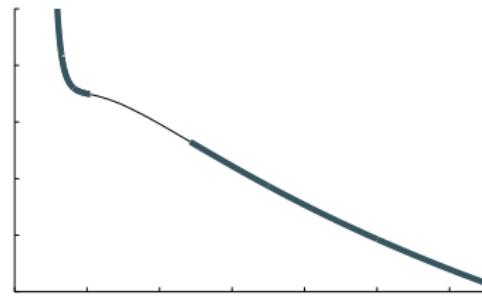
## A Curvature Dependent Pressure Function

Generalised Maxwell construction for  $\tilde{p}^\kappa$

$$\text{Pressure } \tilde{p} = -\frac{d}{d\tau} \tilde{\psi}(\tau)$$



$$\text{Free Energy } \tilde{\psi}$$

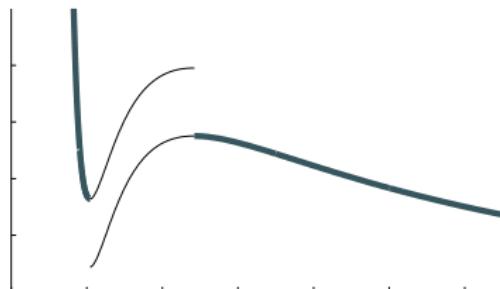


# Microscale

## A Curvature Dependent Pressure Function

Generalised Maxwell construction for  $\tilde{p}^\kappa$

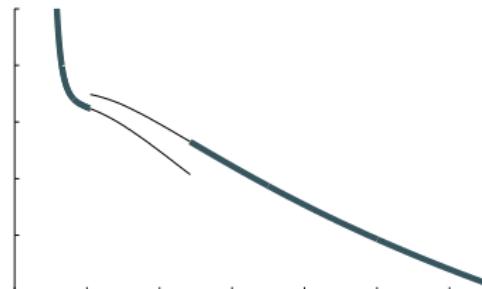
$$\text{Pressure } \tilde{p}^\kappa = -\frac{d}{d\tau} \tilde{\psi}^\kappa(\tau)$$



liquid

$T$   
vapour

$$\text{Free Energy } \tilde{\psi}^\kappa$$



liquid

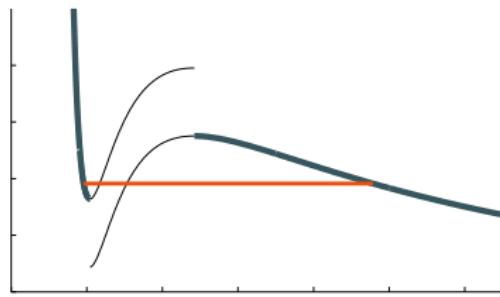
$T$   
vapour

# Microscale

## A Curvature Dependent Pressure Function

Generalised Maxwell construction for  $\tilde{p}^\kappa$

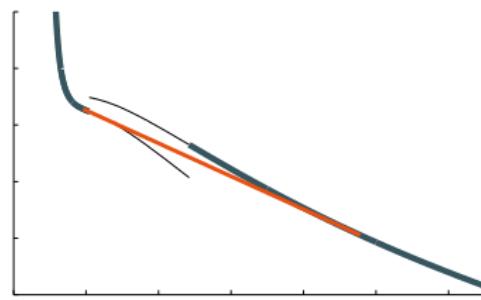
$$\text{Pressure } \tilde{p}^\kappa = -\frac{d}{d\tau} \tilde{\psi}^\kappa(\tau)$$



liquid

$T$   
vapour

$$\text{Free Energy } \tilde{\psi}^\kappa$$



liquid

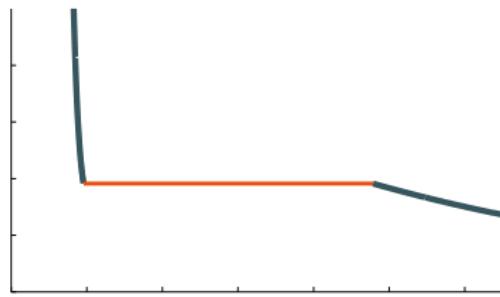
$T$   
vapour

# Microscale

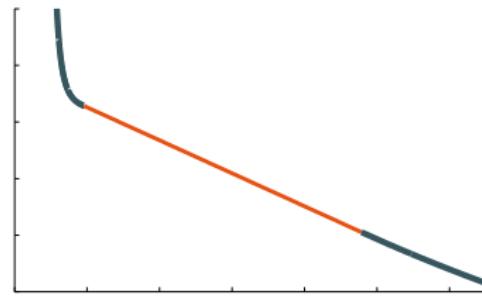
## A Curvature Dependent Pressure Function

Generalised Maxwell construction for  $\tilde{p}^\kappa$

$$\text{Pressure } \tilde{p}^\kappa = -\frac{d}{d\tau} \tilde{\psi}^\kappa(\tau)$$



$$\text{Free Energy } \tilde{\psi}^\kappa$$



# Microscale Solution

$$\begin{pmatrix} \tau \\ u \end{pmatrix}_t + \begin{pmatrix} -u \\ \tilde{p}^\kappa(\tau) \end{pmatrix}_\xi = \begin{pmatrix} 0 \\ 0 \end{pmatrix} \quad \text{in } \mathbb{R} \times (0, \infty). \quad (\text{Euler}^\kappa)$$

## Theorem [Z 2012]

The Riemann problem for  $(\text{Euler}^\kappa)$  with  $|\kappa| < C$  has a unique (Liu) entropy solution. The solution consist of elementary waves and exactly one phase boundary satisfying  $(\text{RH}^\kappa)$ .

- entropy dissipativity condition:

$$\sigma [\![ E^\kappa(\tau, u) ]\!] - [\![ F^\kappa(\tau, u) ]\!] \leq 0,$$

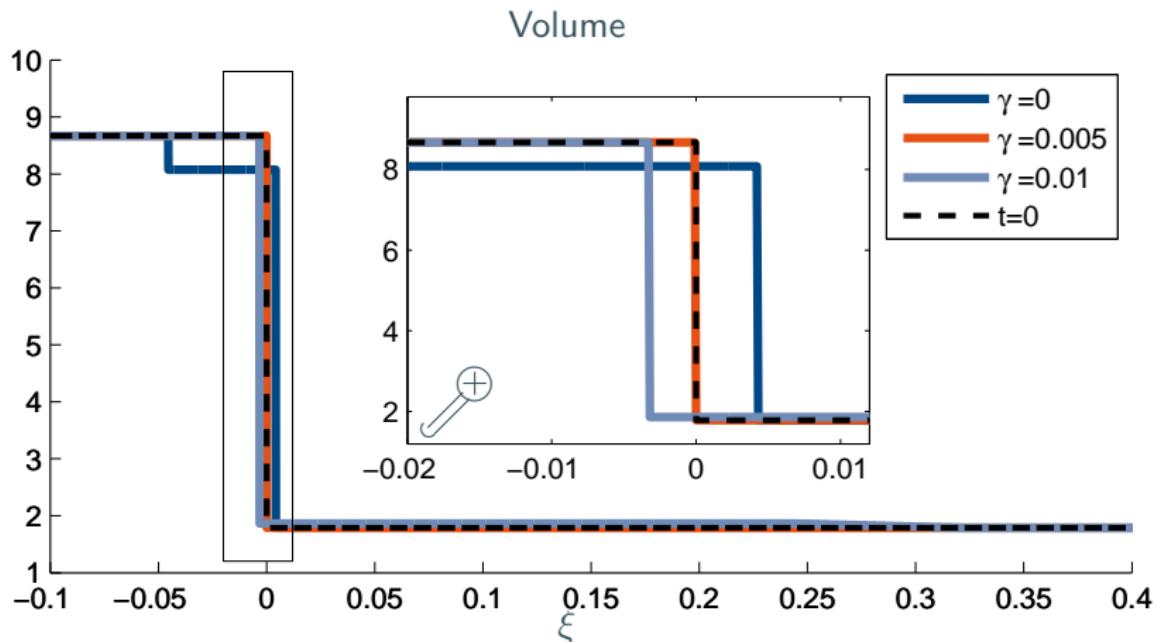
$$E^\kappa(\tau, u) = \frac{1}{2}u^2 + \tilde{\psi}^\kappa(\tau), \quad F^\kappa(\tau, u) = u \tilde{p}^\kappa(\tau).$$

For a construction, we follow [Godlewski, Seguin 2006].

# Microscale

## Solution

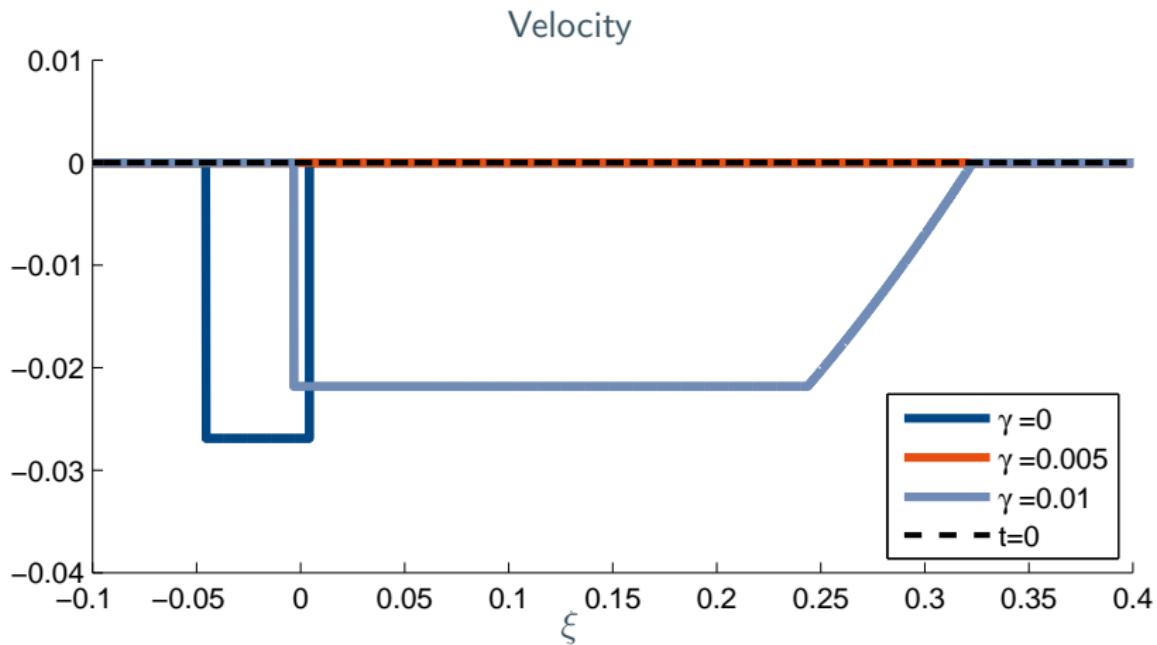
Solution for a bubble of radius 1  
with different surface tension coefficients  $\gamma$ .



# Microscale

## Solution

Solution for a bubble of radius 1  
with different surface tension coefficients  $\gamma$ .



# Application of the HMM in 2d (with P. Engel)

## Macroscale Solver

- Finite Volume method for the 2d Euler system
- unstructured grid
- cut cells along the level set zero
- level set driving

$$\frac{\partial \varphi}{\partial t} + (\sigma \mathbf{n}) \cdot \operatorname{grad} \varphi = \mathbf{0}$$

- curvature

$$\kappa = \operatorname{div} \left( \frac{\operatorname{grad} \varphi}{|\operatorname{grad} \varphi|} \right)$$

# Application of the HMM in 2d

$$\begin{aligned}\varrho(x, 0) &= \\ &\begin{cases} 0.3191 & : \text{inside} \\ 1.8063 & : \text{outside}, \end{cases} \\ v(x, 0) &= \mathbf{0}, \\ \gamma &= 0.001\end{aligned}$$

# Application of the HMM in 2d

$$\begin{aligned}\varrho(x, 0) &= \\ &\begin{cases} 0.3191 & : \text{inside} \\ 1.8063 & : \text{outside}, \end{cases} \\ v(x, 0) &= \mathbf{0}, \\ \gamma &= 0\end{aligned}$$

# Application of the HMM in 2d

## Solution

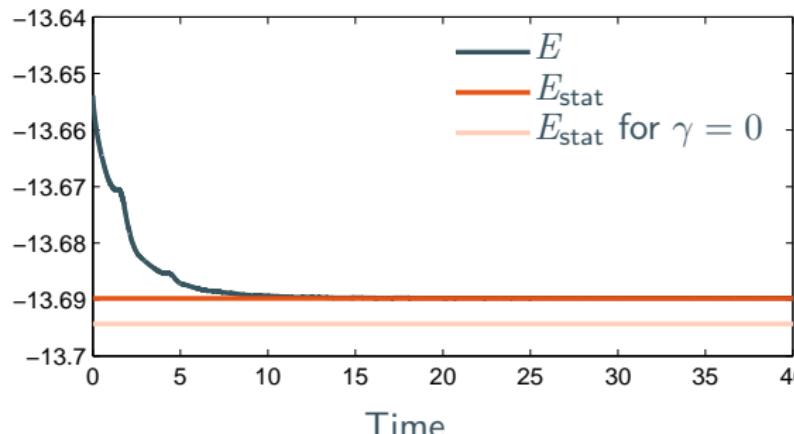
$$\begin{aligned}\varrho(x, 0) &= \\ &\begin{cases} 0.35 & : |x| \leq 0.4 \\ 1.9 & : |x| > 0.4, \end{cases} \\ v(x, 0) &= \mathbf{0}, \\ \gamma &= 0.0025\end{aligned}$$

# Application of the HMM in 2d

**Energy (cf. [Gurtin 1985])**

$$E(\varrho, \mathbf{v}) = \int_{\Omega} \frac{1}{2} \varrho |\mathbf{v}|^2 + W(\varrho) \, dx + \gamma |\Gamma|, \quad W(\varrho) = \varrho \tilde{\psi}(1/\varrho),$$

$$E_{\text{stat}} = \min \left\{ E(\varrho, \mathbf{0}) \mid \int_{\Omega} \varrho \, dx = \text{const.} \right\}$$



# Outlook

## Macroscale

- implementation in Dune / C++
- interface tracking
- curvature reconstruction

## Microscale

- Navier-Stokes-Korteweg model

# Literature

[Engquist et al. 2007] W. E, B. Engquist, X. Li, W. Ren, and E. Vanden-Eijnden.

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