

ENTROPY FORMULATION FOR FORWARD-BACKWARD PARABOLIC EQUATION

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25/06/2012

Collaboration

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Introduction

Forward-backward parabolic equation

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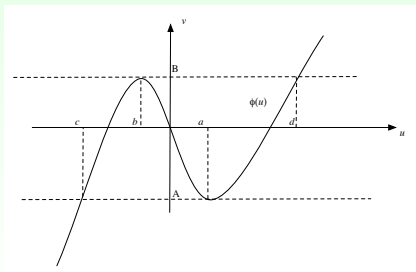
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Example 1

Model of phase separation

$$\phi'(u) > 0 \text{ if } u \in (-\infty, b) \cup (a, \infty), \phi'(u) < 0 \text{ if } u \in (b, a);$$



Example 2

Model of image processing (1D), Perona–Malik equation

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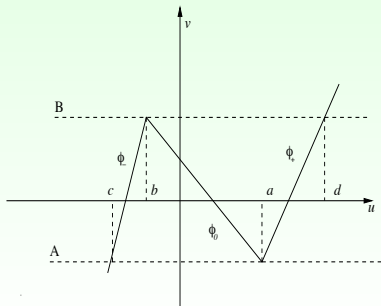
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Hollig (Trans. Amer. Math. Soc. 83) ϕ piecewise linear, there is an infinite number of solutions of the Neumann boundary problem.



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Problem is ill-posed since some relevant physical terms are neglected

Phase transition, Cahn–Hilliard equation

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In the following $\delta = 0$, ϕ is of cubic type.

Novick Cohen-Pego (*Trans. Amer. Math. Soc.* 1991) study the viscosity problem

$$\begin{cases} u_t = \Delta v & \text{in } \Omega \times (0, T] =: Q_T \\ \frac{\partial v}{\partial \nu} = 0 & \text{in } \partial\Omega \times (0, T] \\ u = u_0 & \text{in } \Omega \times \{0\}, \end{cases} \quad (2)$$

where

$$v := \phi(u) + \epsilon u_t \quad (\epsilon > 0), \quad (3)$$

is the *chemical potential*, $\Omega \subseteq \mathbf{R}^n$ is bounded, $\partial\Omega$ regular, $T > 0$.

Partial differential equation in (2) can be rewritten

$$u_t = -\frac{1}{\epsilon}(I - (I - \epsilon\Delta)^{-1})\phi(u)$$

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Moreover $v = (I - \epsilon\Delta)^{-1}\phi(u)$.

Using the standard theory of ODE in the Banach spaces we have

Theorem

(Novick Cohen-Pego) Given $u_0 \in L^\infty(\Omega)$, $\epsilon > 0$ there exists a unique solution (u_ϵ, v_ϵ) defined in $(0, T_\epsilon)$, $u_\epsilon \in C^1([0, T_\epsilon], L^\infty(\Omega))$.

A priori estimates

For every $g \in C^1(\mathbb{R})$ such that $g' \geq 0$

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Then

$$\begin{aligned} [G(u_\epsilon)]_t &= \operatorname{div} \left[g(v_\epsilon) \nabla v_\epsilon \right] - g'(v_\epsilon) |\nabla v_\epsilon|^2 + \\ &\quad - \frac{1}{\epsilon} \left[g(\phi(u_\epsilon)) - g(v_\epsilon) \right] (\phi(u_\epsilon) - v_\epsilon). \end{aligned}$$

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Integrating in Ω and using boundary condition we have

$$\frac{d}{dt} \int_{\Omega} G(u_\epsilon(x, t)) dx \leq 0$$

that gives a priori estimates in L^∞ . Moreover choosing $g(u) \equiv u$ we have

$$\iint_{Q_T} \left\{ |\nabla v^\epsilon|^2 + \epsilon |\partial_t u^\epsilon|^2 \right\} dx dt \leq C_2.$$

Entropy formulation

In analogy with conservation laws we characterize an entropy solution of problem

$$\left\{ \begin{array}{ll} u_t = \Delta \phi(u) & \text{in } \Omega \times (0, T] = Q_T \\ \frac{\partial \phi(u)}{\partial \nu} = 0 & \text{in } \partial\Omega \times (0, T] \\ u = u_0 & \text{in } \Omega \times \{0\}, \end{array} \right. \quad (4)$$

as that obtained as limit of the solutions of problem (2) when $\epsilon \rightarrow 0^+$.

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For every $\epsilon > 0$ and $g \in C^1(\mathbb{R})$, $g' \geq 0$ we have

$$\iint_{Q_T} \left\{ G(u^\epsilon) \psi_t - g(v^\epsilon) \nabla v^\epsilon \cdot \nabla \psi - g'(v^\epsilon) |\nabla v^\epsilon|^2 \psi \right\} \geq 0 \quad (5)$$

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The idea is to pass in the limit in (5) to characterize an entropy solution of (4).

Plotnikov's results

Study of the singular limit, Plotnikov (J. Math. Sci. 1993).

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Using previous a priori estimate we deduce that there exist two subsequence $\{u^{\epsilon_n}\}$, $\{v^{\epsilon_n}\}$ and a couple (u, v) $u \in L^\infty(Q_T)$, $v \in L^\infty(Q_T) \cap L^2((0, T); H^1(\Omega))$ such that for every $T > 0$:

$$u^{\epsilon_n} \xrightarrow{*} u \quad \text{in } L^\infty(Q_T),$$

$$v^{\epsilon_n} \xrightarrow{*} v \quad \text{in } L^\infty(Q_T),$$

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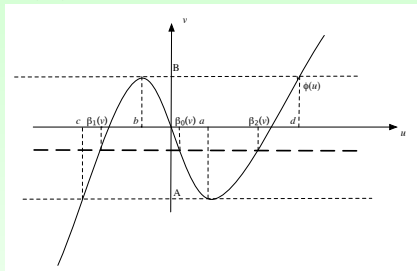
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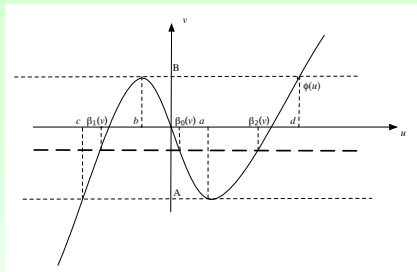
Plotnikov gives a characterization of the Young measure $\nu(x, t)$ associate to the sequence $\{u^{\epsilon_n}\}$ proving that this is given by superposition of Dirac measures,

$$\nu_{(x,t)}(\tau) = \sum_{i=0}^2 \lambda_i(x, t) \delta(\tau - \beta_i(v(x, t))) \quad \text{a.e in } Q_T$$

where $\beta_i(v)$, $i = 0, 1, 2$ are the three branches of the graph $v = \phi(u)$

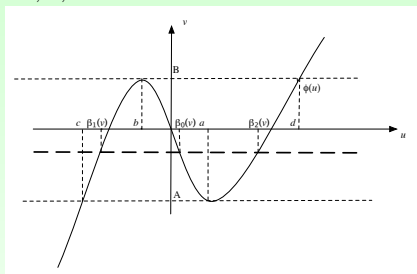


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Moreover, $0 \leq \lambda_i \leq 1$ e $\sum_{i=0}^2 \lambda_i(x, t) = 1$.

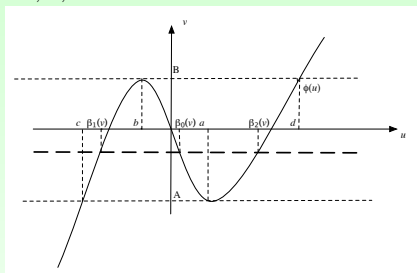
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It could be proved that a subsequence of $\{v^{\epsilon_n}\}$ converges strongly to the function v and (u, v) satisfies equation $u_t = \Delta v$ in the weak sense,

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It could be proved that a subsequence of $\{v^{\epsilon_n}\}$ converges strongly to the function v and (u, v) satisfies equation $u_t = \Delta v$ in the weak sense, meanwhile u satisfies $u_t = \Delta(\phi(u))$ in the sense of measured valued solutions.

Letting $\epsilon_n \rightarrow 0^+$ in the viscous entropy inequality

$$\iint_{Q_T} \left\{ G(u^{\epsilon_n}) \psi_t - g(v^{\epsilon_n}) \nabla v^{\epsilon_n} \cdot \nabla \psi - g'(v^{\epsilon_n}) |\nabla v^{\epsilon_n}|^2 \psi \right\} dx dt +$$

$$\int_{\Omega} G(u_0(x)) \psi(x, 0) dx \geq 0$$

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we have

$$\iint_{Q_T} \left\{ \bar{G}(u) \psi_t - g(v) \nabla v \cdot \nabla \psi - g'(v) |\nabla v|^2 \psi \right\} dx dt +$$

$$\int_{\Omega} G(u_0(x)) \psi(x, 0) dx \geq 0 \tag{6}$$

where $\bar{G}(u) = \sum_{i=0}^2 \lambda_i G(\beta_i(v))$.

Entropy solution

Entropy solution

Definition

Given $u_0 \in L^\infty(\Omega)$ an entropy solution of problem forward–backward (4) is given by the functions $\lambda_i \in L^\infty(Q_T)$, $i = 0, 1, 2$, $u \in L^\infty(Q_T)$, $v \in L^\infty(Q_T) \cap L^2((0, T), H^1(\Omega))$. Such that

(i) $\sum_{i=0}^2 \lambda_i = 1$, $\lambda_i \geq 0$, $u = \sum_{i=0}^2 \lambda_i \beta_i(v)$;

(ii) $u_t = \Delta v$ in the weak sense;

(iii) u and v satisfy (6) for every $g \in C^1(\mathbb{R})$, $g' \geq 0$ (entropy condition).

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Problems: Existence in a stronger sense? Uniqueness? Study of the evolution of the different phases.

Two phase entropy solution

Case $n = 1$. Let $\Omega = (-L, L)$, $u_0 \leq b$ in $(-L, 0)$, $u_0 \geq a$ in $(0, L)$, initial data in the two stable phases.

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an interface ξ , $\xi(0) = 0$,

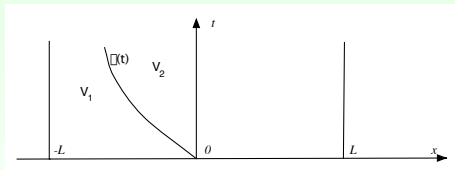
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$V_1 := \{(x, t) \in Q_T \mid -L \leq x < \xi(t), t \in (0, T)\}$, $V_2 := Q_T \setminus \bar{V}_1$



An entropy solution is a triple of functions (ξ, u, v) such that :

(a) $\xi \in C^1((0, T))$, $\xi(0) = 0$, $u \in L^\infty(Q_T)$,

$v \in L^\infty(Q_T) \cap L^2((0, T); H^1((-L, L))$;

(b) u, v satisfy

$$u = \beta_i(v) \text{ in } V_i \quad (i = 1, 2) \quad (v = \phi(u));$$

(c) $u_t = v_{xx}$ in the weak sense, entropy condition, boundary and initial condition.

Determinate conditions for the interface.

A. Terracina

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Theorem

(Evans–Portilheiro *Math. Models Methods Appl. Sci.* (2004))

Let u, v, ξ a two phase entropy “piecewise regular” solution
 then

(i) Rankine-Hugoniot condition:

$$\xi' = -\frac{[v_x]}{[u]} \quad \text{a.e on } \gamma := \{(\xi(t), t) : t \in (0, T)\}.$$

(ii) entropy condition:

$$\xi' [G(u)] \geq -g(v)[v_x] \quad \text{a.e. su } \gamma.$$

where $[h] := h(\xi(t)+, t) - h(\xi(t)-, t)$.

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Condition for the phase change.

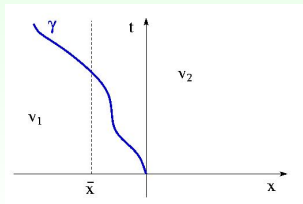
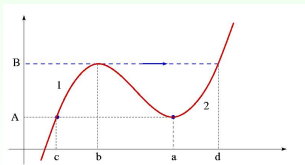
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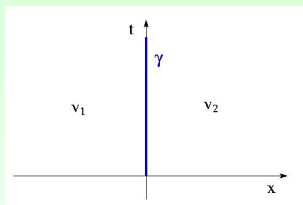
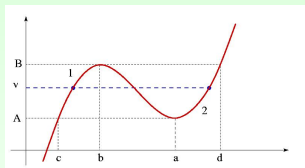
Theorem

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Condition for the phase change. We can pass from phase 1 to phase 2 only if $v = B$



If $v \in (A, B)$ phase does not change

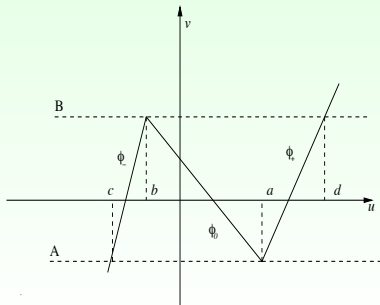


Piecewise linear case

$$\phi(u) = \begin{cases} \phi_-(u) & \text{if } u \leq b \\ \phi_0(u) & \text{if } b < u < a \\ \phi_+(u) & \text{if } u \geq a, \end{cases}$$

where

$$\phi_{\pm}(u) := \alpha_{\pm} u + \beta_{\pm}, \quad \phi_0(u) := \frac{A(u - b) - B(u - a)}{a - b}.$$



Theorem

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i) $\alpha_- u_0(0-) + \beta_- = \alpha_+ u_0(0+) + \beta_+ \in (A, B)$;

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Then there exists $\tau > 0$ such that the two phase problem has solution in $\mathbb{R} \times (0, \tau)$.

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PROOF (idea)

Local existence using two auxiliary problems: moving boundary problem and steady boundary problem.

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We can extend in time the solution until a first time τ such that $\xi'(\tau) = 0$ and $\phi(u(\xi(\tau), \tau)) = A$ or B (critical case).

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It is necessary to study the critical case

Theorem

(T. Siam J. Mat. Anal. 2011) Let (ξ, u) be a solution of the two phase problem in Q_τ . Let $t_1 < \tau$ such that in (t_1, τ) the solution is given by the solution either of the moving boundary problem or of the steady boundary problem. Then there exists $t_2 > \tau$ such that the solution of the two phase problem can be extended in $(0, t_2)$.

Nolinear ϕ

Smarrazzo, T. (Discrete and Continuous dynamical systems A.) Rankine–Hugoniot has to be understood in a weak sense, the same for the entropy condition (theory of vector fields with divergence–measure). We obtain the same compatibility condition for the two phase problem.

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Uniqueness

Theorem

There exists at most one two–phase solution of problem .

Let (u^1, v^1, ξ^1) and (u^2, v^2, ξ^2) be two two–phase solutions then we have

$$\iint_{Q_T} \{(u^1 - u^2)\psi_t - (v_x^1 - v_x^2)\psi_x\} dxdt = 0 \quad (7)$$

for every test function $\psi \in H^1(Q_T) \cap C(\overline{Q_T})$ such that $\psi(\cdot, T) \equiv 0$.

Choosing

$$\psi(x, t) := \int_T^t \{v^1(x, s) - v^2(x, s)\} ds, \quad (8)$$

equation reads

$$\iint_{Q_T} (u^1 - u^2)(v^1 - v^2) dx dt = \iint_{Q_T} (v_x^1 - v_x^2) \left(\int_T^t (v_x^1(x, s) - v_x^2(x, s)) ds \right) dx dt. \quad (9)$$

Concerning the right-hand side of (9) we have

$$\begin{aligned} & \iint_{Q_T} (v_x^1 - v_x^2) \left(\int_T^t (v_x^1(x, s) - v_x^2(x, s)) ds \right) dx dt \quad (10) \\ &= \frac{1}{2} \int_{\omega_1}^{\omega_2} \int_0^T \frac{d}{dt} \left(\left[\int_T^t (v_x^1(x, s) - v_x^2(x, s)) ds \right]^2 \right) dt dx \leq 0. \end{aligned}$$

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Suppose that one of the following conditions is satisfied

- there exists $\delta > 0$ such that $\phi(u_0) < B$ in $(0, \delta)$;
- there exists $\delta > 0$ such that $\phi(u_0) > A$ in $(-\delta, 0)$;

then there exists $\tau > 0$ such that u_{ϵ_n} converges strongly to the solution of the two phase problem with fixed boundary ($\xi \equiv 0$) in $(-L, L) \times (0, \tau)$.

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If $\phi(u_0(\cdot)) \leq B$ in $(-L, x_0)$ and $v^\epsilon(x_0, \cdot) \leq B$ in $(0, t_0)$ then
 $\phi(u^\epsilon), v^\epsilon \leq B$ in $(-L, x_0) \times (0, t_0)$.

Moreover if $u_0(\cdot) \leq b$ $(-L, x_0)$ then $u^\epsilon \leq b$ in $(-L, x_0) \times (0, t_0)$.

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$$\|v_t^{\epsilon_k}\|_{L^2(Q_{\tau_{\epsilon_k}})} + \|v_x^{\epsilon_k}\|_{L^\infty(0, \tau_{\epsilon_k}; L^2((-L, L))} \leq C \quad (11)$$

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Using (11) we can pass to the limit strongly and prove that there exists $\tau > 0$ such that $\tau_{\epsilon_k} \geq \tau$ for k large enough.