

Adaptive measure-valued coupling of non-linear hyperbolic PDEs

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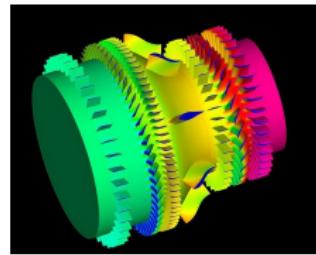
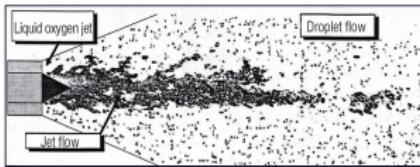
Numerical simulation of systems with multiscale phenomena

- Context:
- complex systems decomposed into sub-systems simulated by specific softwares with different modelings
 - need to couple problems corresponding to distinct modeling levels

Examples: multiphase flows, turbulent flows, porous/non-porous media, kinetic/hydrodynamic representations, etc.

- Aim:
- simulating the whole system with an efficient and reliable coupling method [Godlewski et al. '04 '05, Ambroso et al. '08]
 - coupling hyperbolic systems of conservation laws with non-homogeneous closure laws (i.e. depending on space variables)

- Application:
- modelling with adaptive simplified closure laws
 - application to incomplete EOS [Menikoff & Plohr '89]



Outline

1 Model problem

- Euler equations
- Non-homogeneous equation of state

2 Numerical method

- Coupling procedure
- Relaxation method
- Discretization method

3 Numerical experiments

- Simplified EOS
- Space time adaptive EOS

4 Concluding remarks

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Euler equations with heterogeneity in space

$$\begin{aligned}\partial_t \mathbf{u} + \partial_x \mathbf{f}(\mathbf{u}, x) &= 0, \quad \forall x \in \mathcal{D}(t), \quad t > 0 \\ \mathbf{u}(x, 0) &= \mathbf{u}_0(x), \quad \forall x \in \mathbb{R}\end{aligned}$$

- Time-dependant space domain:

$$\mathcal{D}(t) = \bigcup_{i=1}^{n(t)} (a_{i-1}(t), a_i(t)) \quad \text{with} \quad a_0(t) = -\infty, \quad a_{n(t)}(t) = +\infty, \quad t > 0$$

- Conservative variables:

$$\mathbf{u} = (\rho, \rho u, \rho E)^\top$$

- Heterogeneous physical fluxes in x through pressure definition:

$$\mathbf{f}(\mathbf{u}, x) = \begin{pmatrix} \rho u \\ \rho u^2 + p(\mathbf{u}, x) \\ (\rho E + p(\mathbf{u}, x)) u \end{pmatrix}$$

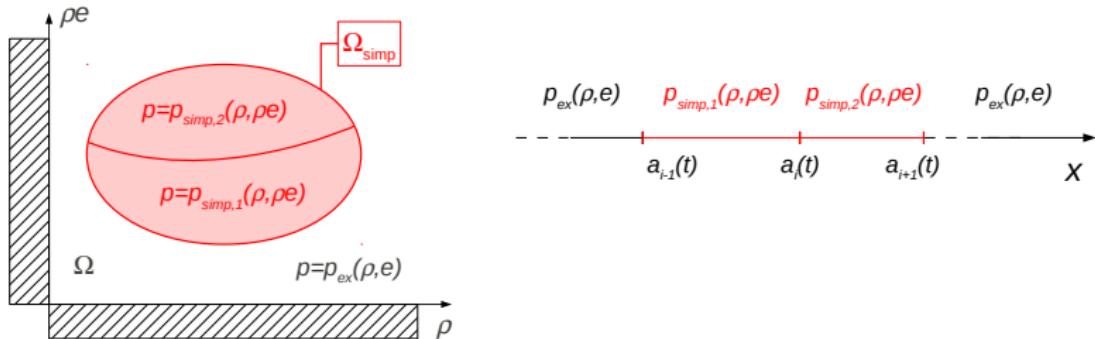
- Set of admissible states:

$$\Omega = \left\{ \mathbf{u} \in \mathbb{R}^3 : \rho > 0, u \in \mathbb{R}, e = E - \frac{u^2}{2} > 0 \right\}$$



Non-homogeneous equation of state (EOS)

$$p(\mathbf{u}, x) = \begin{cases} p_{ex}(\rho, e) & \text{if } \mathbf{u}(x, t) \in \Omega \setminus \Omega_{simp} \\ p_{simp}(\rho, \rho e) & \text{if } \mathbf{u}(x, t) \in \Omega_{simp} \end{cases}$$



- EOS:
- exact (but complex) EOS $p_{ex}(\rho, e)$
 - incomplete (but simple) EOS $p_{simp}(\rho, \rho e)$

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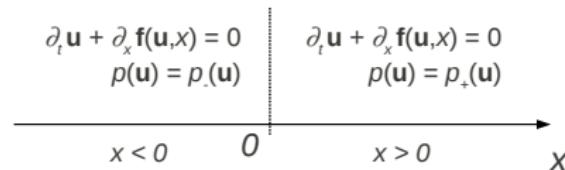
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Coupling procedure

Consider the model problem

$$p(\mathbf{u}, x) = \begin{cases} p_-(\mathbf{u}) & \text{if } x < 0 \\ p_+(\mathbf{u}) & \text{if } x > 0 \end{cases}$$



Coupling prescribed by a vector-valued Dirac measure concentrated at $x = 0$:

$$\begin{aligned} \partial_t \mathbf{u} + \partial_x \mathbf{f}(\mathbf{u}, x) &= \mathcal{M}(\mathbf{u}(0^+, t), \mathbf{u}(0^-, t)) \delta_{x=0}, & \forall x \in \mathbb{R}, \quad t > 0 \\ \mathbf{u}(x, 0) &= \mathbf{u}_0(x), & \forall x \in \mathbb{R} \end{aligned}$$

Coupling procedure (cont.)

Coupling condition

$$\mathcal{M}(\mathbf{u}(0^+, t), \mathbf{u}(0^-, t)) = \mathbf{f}(\mathbf{u}(0^+, t), 0^+) - \mathbf{f}(\mathbf{u}(0^-, t), 0^-), \quad \forall t > 0$$

Example

conservative coupling (continuity of physical fluxes):

$$\mathbf{f}(\mathbf{u}(0^+, t), 0^+) = \mathbf{f}(\mathbf{u}(0^-, t), 0^-)$$

thus

$$\mathcal{M} \equiv 0$$

v-state coupling (continuity of physical variables $\mathbf{v} = \gamma^{-1}(\mathbf{u})$):

$$\gamma_+^{-1}(\mathbf{u}(0^+, t)) = \gamma_-^{-1}(\mathbf{u}(0^-, t))$$

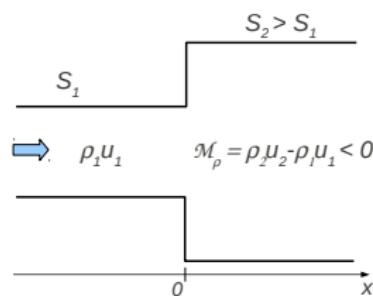
e.g. for $\mathbf{v} = (\rho, u, p)^\top$ we get

$$\mathcal{M} = (0, 0, \rho e(0^+, t)u(0^+, t) - \rho e(0^-, t)u(0^-, t))^\top$$

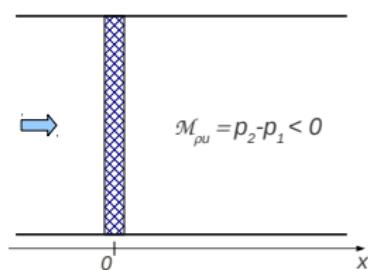
[Godlewski et al. '05, Ambroso et al. '08a. '07, Galie '09]

Non-conservative framework

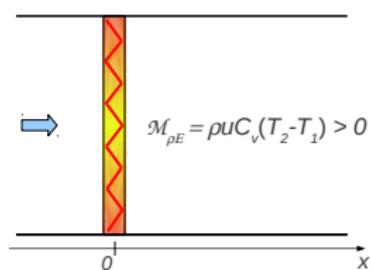
$$\mathcal{M} = \begin{pmatrix} \mathcal{M}_\rho \\ \mathcal{M}_{\rho u} \\ \mathcal{M}_{\rho E} \end{pmatrix}$$



(a) sudden pipe expansion



(b) pressure loss



(c) heat transfert

Relaxation method

- Approximate the Euler eq. with a Suliciu's relaxation procedure ($\lambda \gg 1$):

$$\begin{aligned}\partial_t \mathbf{U} + \partial_x \mathbf{F}(\mathbf{U}, x) &= \mathcal{M} \delta_{x=0} + \lambda \mathcal{R}(\mathbf{U}, x), \quad \forall x \in \mathbb{R}, \quad t > 0 \\ \mathbf{U}(x, 0) &= \mathbf{U}_0(x), \quad \forall x \in \mathbb{R}\end{aligned}$$

with

$$\mathbf{U} = (\rho, \rho u, \rho E, \rho \pi)^\top$$

and

$$\mathbf{F}(\mathbf{U}, x) = \begin{pmatrix} \rho u \\ \rho u^2 + \pi \\ (\rho E + \pi)u \\ (\rho \pi + a^2)u \end{pmatrix}$$

and

$$\mathcal{R}(\mathbf{U}, x) = (0, 0, 0, \rho(p(u, x) - \pi))^\top$$

- Advantage:** the system has 4 LD fields $\{u - a\tau, u, u, u + a\tau\}$
- Whitham condition:

$$a^2 > \max\{-(\partial_\tau p)_s : \text{for all } \tau, s \text{ under consideration}\}$$

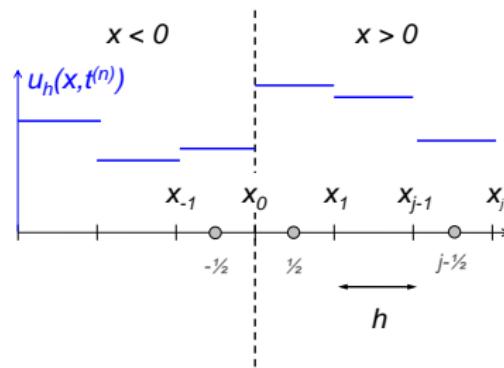
[Jin & Xin '95, Suliciu '98, Ambroso et al. '08b]

Discretization method

Numerical solution

- Let $\bigcup_{j \in \mathbb{Z}} [x_{j-1}, x_j)$ a partition of the domain \mathbb{R} with $x_j = jh$ and $h > 0$
- Look for a piecewise constant solution \mathbf{U}_h at time step $t^{(n)} = n\Delta t$:

$$\mathbf{U}_h(x, t^{(n)}) = \mathbf{U}_{j-\frac{1}{2}}^{(n)}, \quad \forall x \in (x_{j-1}, x_j), \quad n > 0$$



Discretization method away from the coupling interface

2-step algorithm from $t^{(n)}$ to $t^{(n+1)}$

- ❶ solve the relaxation system with $\lambda = 0$:

- ▶ Euler method for the time discretization
- ▶ Finite volume method for the space discretization:

$$\mathbf{U}_{j-\frac{1}{2}}^{(n+1)} = \mathbf{U}_{j-\frac{1}{2}}^{(n)} - \frac{\Delta t}{h} \left(\mathbf{g}_- (\mathbf{U}_{j-\frac{1}{2}}^{(n)}, \mathbf{U}_{j+\frac{1}{2}}^{(n)}) - \mathbf{g}_- (\mathbf{U}_{j-\frac{3}{2}}^{(n)}, \mathbf{U}_{j-\frac{1}{2}}^{(n)}) \right), \quad j < 0$$

$$\mathbf{U}_{j+\frac{1}{2}}^{(n+1)} = \mathbf{U}_{j+\frac{1}{2}}^{(n)} - \frac{\Delta t}{h} \left(\mathbf{g}_+ (\mathbf{U}_{j+\frac{1}{2}}^{(n)}, \mathbf{U}_{j+\frac{3}{2}}^{(n)}) - \mathbf{g}_+ (\mathbf{U}_{j-\frac{1}{2}}^{(n)}, \mathbf{U}_{j+\frac{1}{2}}^{(n)}) \right), \quad j > 0$$

- ▶ Godunov method:

$$\mathbf{g}_\pm(\cdot, \cdot) = \mathbf{F}^\pm(\mathcal{W}(0, \cdot, \cdot))$$

with \mathcal{W} solution of a local Riemann problem

- ❷ projection of the solution onto the equilibrium state for $\lambda \rightarrow \infty$:

$$\lim_{\lambda \rightarrow \infty} \pi^{(n+1)} - p^{(n+1)} = 0$$

Discretization method at a coupling interface

We consider $\mathcal{W}\left(\frac{x}{t-t^{(n)}}; \mathbf{U}_{-\frac{1}{2}}^{(n)}, \mathbf{U}_{+\frac{1}{2}}^{(n)}, \mathcal{M}^{(n)}\right)$ solution of the Riemann problem:

$$\partial_t \mathbf{U} + \partial_x \mathbf{F}(\mathbf{U}, x) = \mathcal{M}^{(n)}, \quad \forall x \in (x_{-\frac{1}{2}}, x_{+\frac{1}{2}}), \quad t \in (t^{(n)}, t^{(n+1)})$$

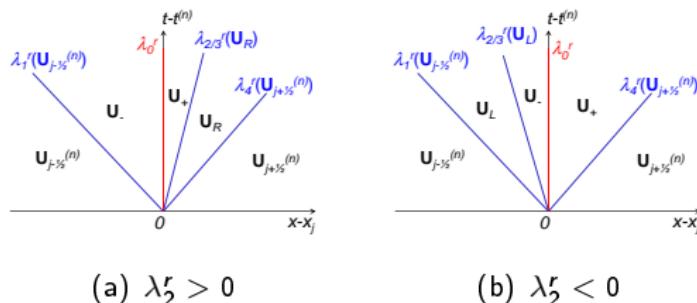
$$\mathbf{U}(x, t^{(n)}) = \begin{cases} \mathbf{U}_{-\frac{1}{2}}^{(n)} & x < 0 \\ \mathbf{U}_{+\frac{1}{2}}^{(n)} & x > 0 \end{cases}$$

Coupling condition at discrete level:

$$\mathcal{M}^{(n)} = \mathbf{g}_+(\mathbf{U}_{-\frac{1}{2}}^{(n)}, \mathbf{U}_{\frac{1}{2}}^{(n)}) - \mathbf{g}_-(\mathbf{U}_{-\frac{1}{2}}^{(n)}, \mathbf{U}_{\frac{1}{2}}^{(n)})$$

Discretization method at a coupling interface (cont.)

Structure of the solution:



Galié et al. (2009) proved existence and unicity of self-similar solution for $p = p(\rho)$ and $\mathcal{M}^{(n)}$ in a suitable domain with $\mathcal{M}_\rho = 0$.

Extended to $\mathcal{M}_\rho \neq 0$ in the present study

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Shock tube problem

- Riemann problem with shock tube conditions:

$$(\rho, u, p)^\top = \begin{cases} (0.4, 1, 1)^\top, & \text{for } x \leq 0 \\ (0.8, -0.5, 1)^\top, & \text{for } x > 0 \end{cases}$$

- for a gas with an (exact) van der Waals EOS with constant heats

$$(p_{ex} + a\rho^2)(1 - b\rho) = \rho r T$$

with $a = 3$, $b = \frac{1}{3}$, $\frac{r}{C_v} = 0.329$

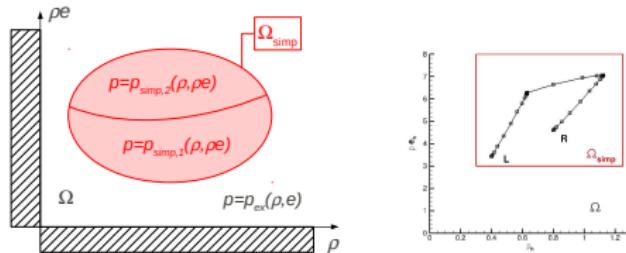
- The simplified EOS is a stiffened equation of Grüneisen type:

$$p_{simp}(\rho, \rho e) = (\gamma - 1)\rho e + c_{ref}^2(\rho - \rho_{ref})$$

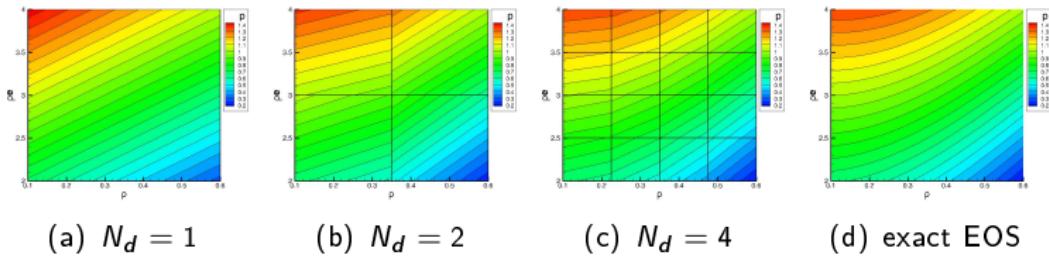
[Calen '88, Menikoff & Plohr '89, Quartapelle et al. '03]

Simplified EOS via diadic hierarchy of Ω_{simp}

- The domain Ω_{simp} is defined s.t. the simplified EOS is valid everywhere:



- Use diadic partition of Ω_{simp} into N_d^2 subdomains:



(a) $N_d = 1$

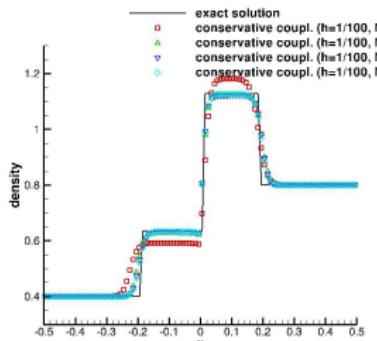
(b) $N_d = 2$

(c) $N_d = 4$

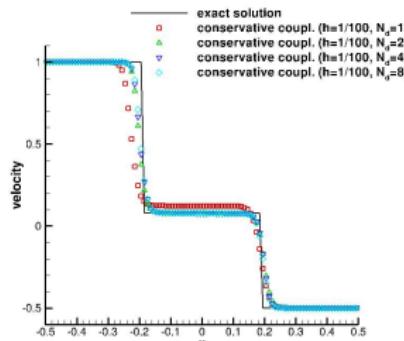
(d) exact EOS

Conservative coupling ($\mathcal{M} \equiv 0$)

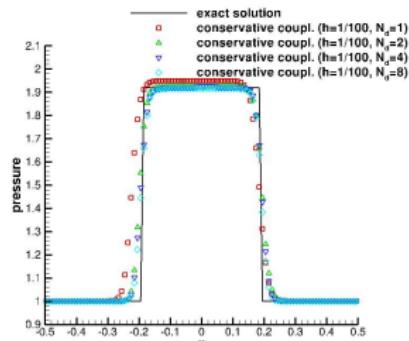
N_d^2 subdomains in Ω_{simp} , $t = 0.13$ and $h = 1/100$



(a) density



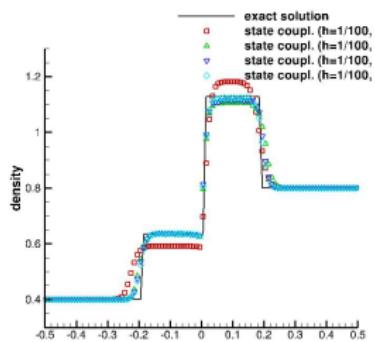
(b) velocity



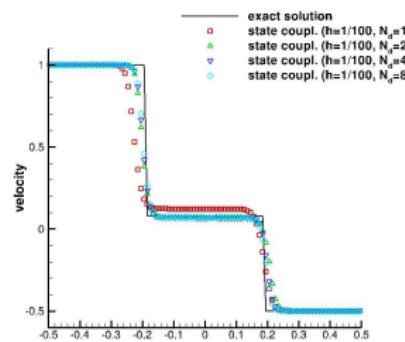
(c) pressure

(ρ, u, p) -state coupling

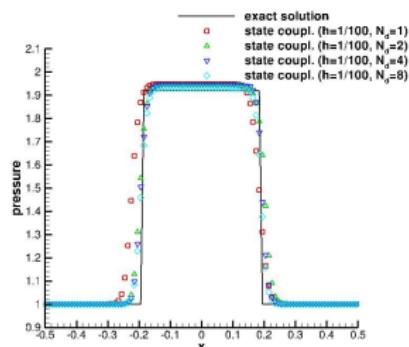
N_d^2 subdomains in Ω_{simp} , $t = 0.13$ and $h = 1/100$



(a) density



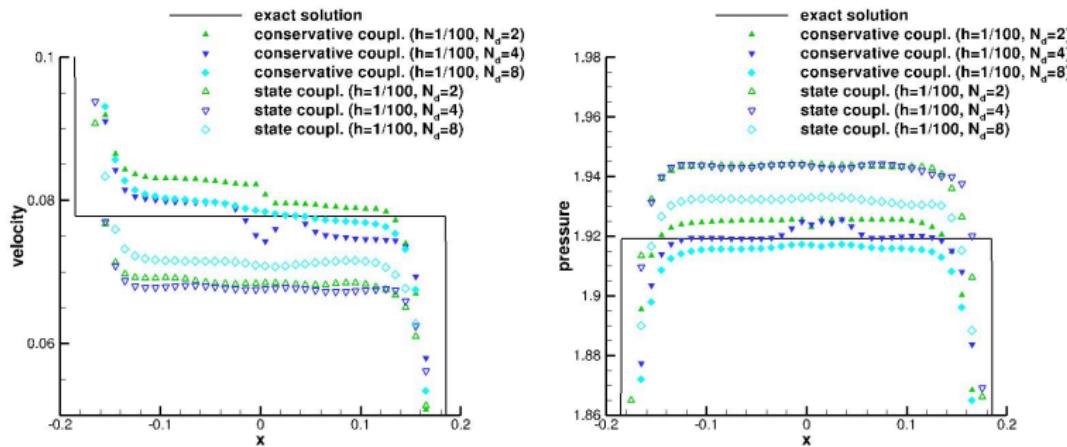
(b) velocity



(c) pressure

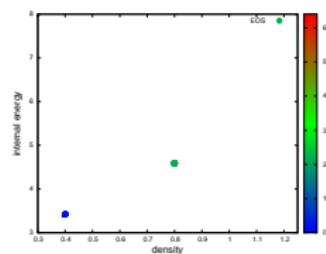
Comparison of coupling methods

N_d^2 subdomains in Ω_{simp} , $t = 0.13$ and $h = 1/100$

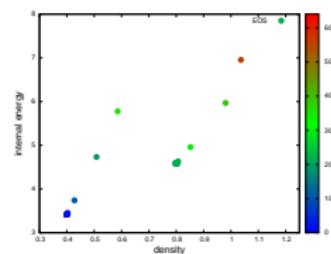


Space time adaptive EOS

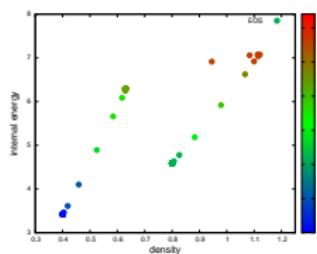
Conservative coupling ($\mathcal{M} \equiv 0$): computation with $N_d^2 = 64$ different EOS ($h = 1/100$)



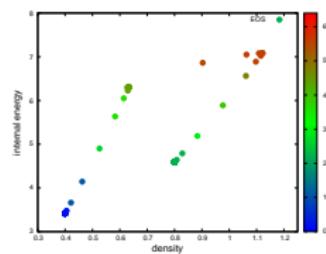
(a) $t = 0$



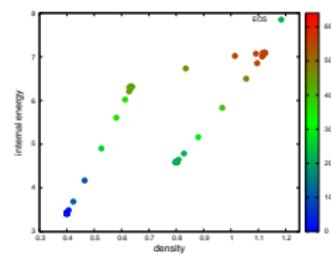
(b) $t = 0.01$



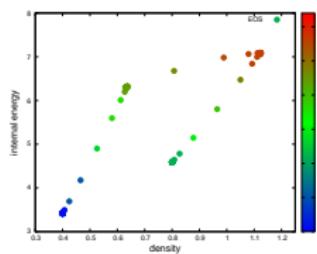
(c) $t = 0.05$



(d) $t = 0.07$



(e) $t = 0.11$



(f) $t = 0.13$

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Adaptive coupling allows to:

- use simplified EOS instead of real gas EOS with complex thermodynamic laws
- reproduce the correct behaviour of real gas dynamics
- use different coupling methods (source term measure offers a general framework)

Future work:

- use mixed coupling methods according to physical requirements
- minimization of convex non-linear cost functions built on the mass \mathcal{M} at the interface

Thank you for your attention

Bibliography

- [1] A. Ambroso, C. Chalons, F. Coquel, E. Godlewski, F. Lagoutière, P.-A. Raviart and N. Seguin, **Coupling of general Lagrangian systems**, *Math. Comp.*, 77 (2008), 909–941.
- [2] A. Ambroso, C. Chalons, F. Coquel, E. Godlewski and P.-A. Raviart, **Relaxation methods and coupling procedures**, *Int. J. Numer. Meth. Fluids*, 56 (2008), 1123–1129.
- [3] A. Ambroso, J.-M. Hérard and O. Hurisse, **A Method to Couple HEM and HRM Two-Phase Flow Models**, *Computers and Fluids*, 38 (2009), 738–756.
- [4] B. Boutin, **Etude mathématique et numérique d'équations hyperboliques non-linéaires : couplage de modèles et chocs non classiques**, Thèse de doctorat, Univ. Pierre et Marie Curie, Paris 6, Nov. 2009.
- [5] H. B. Callen, **Thermodynamics and an introduction to thermostatics**, second edition, John Wiley & Sons, 1988.
- [6] T. Galié, **Couplage interfacial de modèles en dynamique des fluides. Application aux écoulements diphasiques**, Thèse de doctorat, Univ. Pierre et Marie Curie, Paris 6, Mars 2009.
- [7] E. Godlewski and P.-A. Raviart, **The numerical Interface Coupling of Nonlinear Hyperbolic Systems of Conservation Laws. I. The Scalar Case**, *Numer. Math.*, 97 (2004), 81–130.
- [8] E. Godlewski, K.-C. Le Thanh and P.-A. Raviart, **The numerical Interface Coupling of Nonlinear Hyperbolic Systems of Conservation Laws. II. The case of systems**, *M2AN*, 39 (2005), 649–6920.
- [9] S. Jin and Z. Xin, **The Relaxation Schemes for Systems of Conservation Laws in Arbitrary Space Dimension**, *Comm. Pure. Appl. Math.*, 48 (1995), pp. 235–276.
- [10] A. Murrone and P. Villedieu, **Numerical modeling of dispersed two-phase flows**, *The Onera Journal Aerospace Lab*, 2 (2011).
- [11] R. Menikoff and B. J. Plohr, **The Riemann problem for fluid flow of real materials**, *Rev. Mod. Phys.*, 61 (1989), 75–130.
- [12] L. Quartapelle, L. Castelletti, A. Guardone and G. Quaranta, **Solution of the Riemann problem of classical gasdynamics**, *J. Comput. Phys.*, 190 (2003), 118–140.
- [13] J. U. Schlüter, X. Wu, E. V. D. Weide, S. Hahn, J. Alonso and H. Pitsch, **Multi-Code Simulations : a Generalized Coupling Approach**, *AIAA paper 2005-4997*, 2005.
- [14] I. Suliciu, **On the thermodynamics of fluids with relaxation and phase transitions**, *Fluids with relaxation*, *Int. Sci.*, 36 (1998), pp. 921–947.