

A posteriori estimates from approximate solutions of the Euler or Navier-Stokes equations

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(joint work with Carlo Morosi, Politecnico di Milano, carlo.morosi@polimi.it)

Introduction

- I consider the incompressible Euler or Navier-Stokes (NS) Cauchy problem

$$\frac{\partial u}{\partial t} = \nu \Delta u + \mathcal{P}(u, u) + f, \quad \operatorname{div} u = 0, \quad u(x, 0) = u_0(x)$$

$u = u(x, t)$; $x = (x_1, \dots, x_d) \in \mathbf{T}^d := (\mathbf{R}/2\pi\mathbf{Z})^d$ (space-periodic bdy cond.);
 Δ the Laplacian; \mathcal{P} the bilinear map such that, for $v, w : \mathbf{T}^d \rightarrow \mathbf{R}^d$,

$$\mathcal{P}(v, w) := -\mathcal{L}(v \cdot \operatorname{grad} w)$$

with \mathcal{L} the Leray projection onto the space of divergence free vector fields;
 $f = f(x, t)$ the (Leray projected) density of external forces;
Viscosity $\nu = 0$ in the Euler case; $\nu \in (0, +\infty)$ in the NS case.

- The aim is to derive fully quantitative estimates on the exact solution u analyzing *a posteriori* any approximate solution u_a (in particular: estimates on the interval of existence of u , that could be $[0, +\infty)$, and on the distance between u and u_a).
- Concrete examples will be given in space dimension $d = 3$.

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Functional setting on the torus $\mathbf{T}^d := (\mathbf{R}/2\pi\mathbf{Z})^d$

- We use the Fourier basis $e_k(x) := (2\pi)^{-d/2} e^{ik \cdot x}$ ($x \in \mathbf{T}^d, k \in \mathbf{Z}^d$).
- A vector field v on \mathbf{T}^d means a vector distribution $v \in \mathcal{D}'(\mathbf{T}^d, \mathbf{R}^d) \equiv \mathcal{D}'$. This has Fourier coefficients $v_k = \overline{v_{-k}} \in \mathbf{C}^d$, such that $v = \sum_{k \in \mathbf{Z}^d} v_k e_k$.
- The space of zero mean vector fields is

$$\mathcal{D}'_0 := \{v \in \mathcal{D}' \mid \int_{\mathbf{T}^d} v \, dx = 0\} = \{v \in \mathcal{D}' \mid v_0 = 0\};$$

these have Fourier coefficients labelled by $\mathbf{Z}'_0 := \mathbf{Z}^d \setminus \{0\}$.

- The space of divergence free (or solenoidal) vector fields and the Leray projection are

$$\mathcal{D}'_{\Sigma} := \{v \in \mathcal{D}' \mid \operatorname{div} v = 0\} = \{v \in \mathcal{D}' \mid k \cdot v_k = 0 \text{ for } k \in \mathbf{Z}'_0\};$$

$$\mathcal{L} : \mathcal{D}' \rightarrow \mathcal{D}'_{\Sigma}, \quad v \mapsto \mathcal{L}v \text{ such that } (\mathcal{L}v)_k = \mathcal{L}_k v_k \text{ for } k \in \mathbf{Z}'_0,$$

\mathcal{L}_k the projection of \mathbf{C}^d onto the orthogonal complement k^\perp .

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Functional setting on the torus $\mathbf{T}^d := (\mathbf{R}/2\pi\mathbf{Z})^d$

- For $n \in \mathbf{R}$, we consider the n -th Sobolev space of zero mean vector fields

$$\mathbb{H}_0^n := \{v \in \mathbb{D}'_0 \mid \sqrt{-\Delta}^n v \in \mathbb{L}^2\};$$

its inner product and norm are

$$\langle v|w \rangle_n := \langle \sqrt{-\Delta}^n v | \sqrt{-\Delta}^n w \rangle_{L^2} = \sum_{k \in \mathbf{Z}_0^d} |k|^{2n} \overline{v_k} \cdot w_k, \quad \|v\|_n := \sqrt{\langle v|v \rangle_n}.$$

- We also consider the solenoidal subspace

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Some inequalities for $\mathcal{P}(v, w) := -\mathfrak{L}(v \cdot \text{grad} w)$ (on \mathbf{T}^d).

- **Basic inequality:** $\|\mathcal{P}(v, w)\|_n \leq K_n \|v\|_n \|w\|_{n+1}$ for $v \in \mathbb{H}_{\Sigma_0}^n$, $w \in \mathbb{H}_{\Sigma_0}^{n+1}$ and $n > d/2$ (for a suitable constant $K_n > 0$).
- **Kato inequality:** $|\langle \mathcal{P}(v, w) | w \rangle_n| \leq G_n \|v\|_n \|w\|_n^2$ for $v \in \mathbb{H}_{\Sigma_0}^n$, $w \in \mathbb{H}_{\Sigma_0}^{n+1}$ and $n > d/2 + 1$ (for a suitable constant $G_n > 0$).
- **Proposition** (Morosi and P: arXiv:1007.4412v2 [math.AP], 2010; Commun. Pure Appl. Anal., 2012). One can take

$$K_n = \sqrt{\sup_{k \in \mathbb{Z}_0^d} \frac{|k|^{2n}}{(2\pi)^d} \mathcal{K}_n(k)}, \quad \mathcal{K}_n(k) := \sum_{h \in \mathbb{Z}_0^d, h \neq k} \frac{|h \wedge k|^2}{|h|^{2n+2} |k-h|^{2n+2}};$$

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($h \wedge k$ the exterior product; it is the usual vector product, if $d = 3$).

This implies, e.g., that one can take

$$K_2 = 0.335, \quad K_3 = 0.323, \quad K_4 = 0.441, \quad K_5 = 0.510, \quad K_{10} = 2.88;$$

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Cauchy problem for Euler/NS equations (on \mathbf{T}^d).

- Choose $n > d/2 + 1$, $\nu \geq 0$, a forcing $f \in C([0, +\infty), \mathbb{H}_{\Sigma_0}^n)$, and a datum $u_0 \in \mathbb{H}_{\Sigma_0}^{n+2}$. The corresponding Cauchy problem is:

Find $u \in C([0, T), \mathbb{H}_{\Sigma_0}^{n+2}) \cap C^1([0, T), \mathbb{H}_{\Sigma_0}^n)$ such that

$$\frac{du}{dt} = \nu \Delta u + \mathcal{P}(u, u) + f, \quad u(0) = u_0.$$

Here are some basic facts (T. Kato, Lecture Notes in Mathematics **448**,1975; J. T.Beale, T. Kato, A. Majda, Commun. Math. Phys.,1984):

- There is a unique maximal ($:=$ unextendable) solution u of domain $[0, T)$, with $T = T(u_0) \in (0, +\infty]$.
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The theory of approximate solutions. A basic lemma

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$$\frac{d^+ \|w\|_n}{dt}(t_0) \leq \left\| \frac{dw}{dt}(t_0) \right\|_n \stackrel{(1)}{=} \|\text{err}(u_a)(t_0)\| \leq \epsilon_n(t_0),$$

and again we have the desired inequality for $d^+ \|w\|_n / dt$.



The theory of approximate solutions. Control inequalities

- Again, we consider for the Euler/NS Cauchy problem:
 - the maximal solution $u \in C([0, T], \mathbb{H}_{\Sigma_0}^{n+2}) \cap C^1([0, T], \mathbb{H}_{\Sigma_0}^n)$;
 - an approximate solution $u_a \in C([0, T_a], \mathbb{H}_{\Sigma_0}^{n+2}) \cap C^1([0, T_a], \mathbb{H}_{\Sigma_0}^n)$ with differential error, datum error and growth estimators $\epsilon_n, \delta_n, \mathcal{D}_n, \mathcal{D}_{n+1}$.
- **Proposition** (Morosi and P, *Nonlinear Analysis*, 2012). Consider a function $\mathcal{R}_n \in C([0, T_c], [0, +\infty))$, with $T_c \in (0, T_a]$; assume this fulfills the *control inequalities*

$$\frac{d^+ \mathcal{R}_n}{dt} \geq -\nu \mathcal{R}_n + (G_n \mathcal{D}_n + K_n \mathcal{D}_{n+1}) \mathcal{R}_n + G_n \mathcal{R}_n^2 + \epsilon_n, \quad \mathcal{R}_n(0) \geq \delta_n.$$

Then, u and its the existence time T fulfill

$$T \geq T_c; \quad \|u(t) - u_a(t)\|_n \leq \mathcal{R}_n(t) \quad \text{for } 0 \leq t < T_c.$$

In particular, if the control inequalities have a **global solution** \mathcal{R}_n ($T_c = +\infty$), then u is **global as well** ($T = +\infty$) .

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Proof. i) Again, we write $w := u - u_a$. The previous Lemma for $\|w\|_n$ and the control inequalities for \mathcal{R}_n ensure that

$$\frac{d^+ \|w\|_n}{dt}(t) \leq \psi(\|w(t)\|_n, t) \quad \text{for } 0 \leq t < \min(T_a, T), \quad \mathcal{W}_n(0) \leq \delta_n ;$$

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By some standard [comparison lemma](#) *à la* Čaplygin-Laksmikantham, this implies

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and, passing to $u = u_a + w$,

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- The simplest way to solve the control inequalities is to search for a function $\mathcal{R}_n \in C^1([0, T_c), [0, +\infty))$ that satisfies them as equalities:

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- I recall that $e_k(x) := (2\pi)^{-d/2} e^{ik \cdot x}$ for all $k \in \mathbf{Z}^d$. We choose a set of modes

$$G \subset \mathbf{Z}_0^d, \quad G \text{ finite and such that } k \in G \Leftrightarrow -k \in G;$$

we consider the Galerkin subspace and the associated projection

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- The datum error of u_G is $u_0 - u_G(0) = (1 - E^G)u_0 = \sum_{k \in \mathbb{Z}_0^d \setminus G} u_{0k}$;
- The differential error $err(u_G) := du_G/dt - \nu \Delta u_G - \mathcal{P}(u_G, u_G)$ is given by

$$err(u_G) = -(1 - E^G)\mathcal{P}(u_G, u_G) - (1 - E^G)f ;$$

$$(1 - E^G)f = \sum_{k \in \mathbb{Z}_0^d \setminus G} f_k e_k ; \quad (1 - E^G)\mathcal{P}(u_G, u_G) = \sum_{k \in dG} p_k e_k ,$$

$$dG := (G + G) \setminus (G \cup \{0\}) , \quad p_k := -\frac{i}{(2\pi)^{d/2}} \sum_{h \in G} [\gamma_h \cdot (k - h)] \mathfrak{L}_k \gamma_{k-h} .$$

$(G + G := \{p + q \mid p, q \in G\})$. Recall that \mathfrak{L}_k projects onto k^\perp and γ_h are the **Fourier coefficients of u_G** , for $h \in G$; $\gamma_h := 0$ if $h \in \mathbb{Z}_0^d \setminus G$.)

- In order to apply the **control inequalities**, one can **compute directly** the norms

$$\delta_n := \|u_0 - u_G(0)\|_n$$

and, **after solving the Galerkin equations for the coefficients $\gamma_k(t)$,**

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An application. Galerkin approximate solutions

- The datum error of u_G is $u_0 - u_G(0) = (1 - E^G)u_0 = \sum_{k \in \mathbf{Z}_0^d \setminus G} u_{0k} e_k$;
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$$dG := (G + G) \setminus (G \cup \{0\}) , \quad p_k := -\frac{i}{(2\pi)^{d/2}} \sum_{h \in G} [\gamma_h \cdot (k - h)] \mathfrak{L}_k \gamma_{k-h} .$$

$(G + G := \{p + q \mid p, q \in G\})$. Recall that \mathfrak{L}_k projects onto k^\perp and γ_h are the **Fourier coefficients of u_G** , for $h \in G$; $\gamma_h := 0$ if $h \in \mathbf{Z}_0^d \setminus G$.)

- In order to apply the **control inequalities**, one can **compute directly** the norms

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Example. Galerkin method for the BNW datum

- The space dimension is $d = 3$. There is no external forcing: $f = 0$.
- The initial datum of Behr, Nečas and Wu (from now on: BNW datum) is

$$u_0(x_1, x_2, x_3) :=$$

$$2(\cos(x_1+x_2)+\cos(x_1+x_3), -\cos(x_1+x_2)+\cos(x_2+x_3), -\cos(x_1+x_3)-\cos(x_2+x_3))$$

In terms of Fourier modes:

$$u_0 = \sum_{k=\pm a, \pm b, \pm c} u_{0k} e_k, \quad a := (1, 1, 0), \quad b := (1, 0, 1), \quad c := (0, 1, 1);$$

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- In [Morosi and P, *Nonlinear Analysis*, 2012] the Galerkin solution u_G has been computed numerically for the BNW datum and several values of ν , choosing a set of 150 modes

$$G := \{\pm(1, 1, 0), \pm(1, 0, 1), \pm(0, 1, 1), \dots, \pm(3, 3, -2), \pm(3, 3, 2)\}.$$

- Then, the theory of approximate solutions has been applied to u_G in the framework of the spaces $\mathbb{H}_{\Sigma_0}^n, \mathbb{H}_{\Sigma_0}^{n+1}, \mathbb{H}_{\Sigma_0}^{n+2}$ with $n = 3$.
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- The final step has been the analysis of the control Cauchy problem

$$\frac{d\mathcal{R}_3}{dt} = -\nu\mathcal{R}_3 + (G_3\mathcal{D}_3 + K_3\mathcal{D}_4)\mathcal{R}_3 + G_3\mathcal{R}_3^2 + \epsilon_3, \quad \mathcal{R}_3(0) = 0.$$

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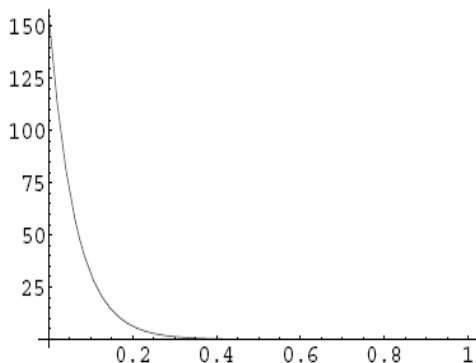
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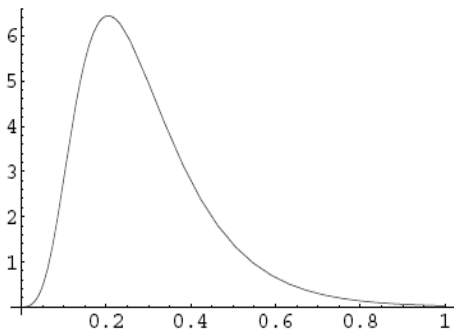
Example. Galerkin method for the BNW datum

Here are some results from the numerical computation of u_G and \mathcal{R}_3 .



$\nu = 8$. Graph of $\mathcal{D}_3(t) = \|u_G(t)\|_3$.

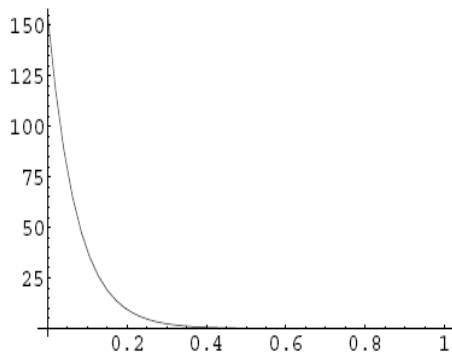
Example. Galerkin method for the BNW datum



$\nu = 8$. Graph of $\mathcal{R}_3(t)$, that appears to be global.

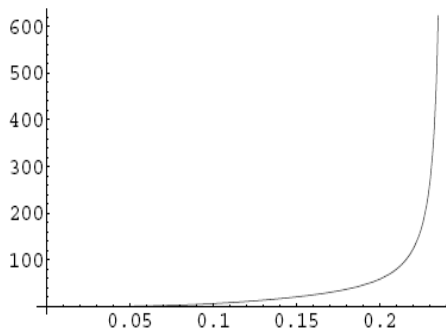
$\|u(t) - u_G(t)\|_3 \leq \mathcal{R}_3(t)$ for all $t \in [0, +\infty)$.

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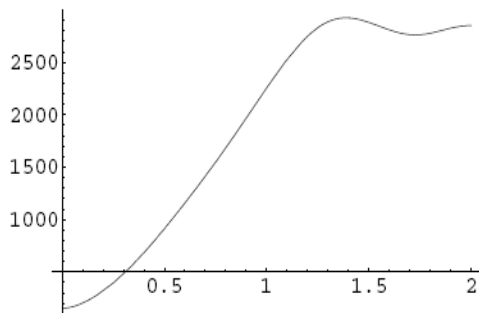
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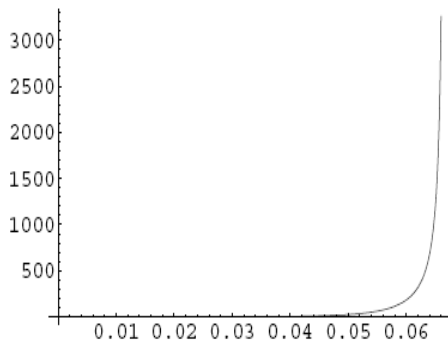
$\nu = 7$. Graph of $\mathcal{R}_3(t)$, that appears to diverge for $t \rightarrow T_c = 0.2386\dots$
 $\|u(t) - u_G(t)\|_3 \leq \mathcal{R}_3(t)$ for $t \in [0, T_c)$.

Example. Galerkin method for the BNW datum



$\nu = 0$. Graph of $\mathcal{D}_3(t) = \|u_G(t)\|_3$.

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Other approaches to build approximate solutions

- Euler or NS equations have **formal power series solutions**

$$u(t) = \sum_{j=0}^{+\infty} u_j t^j \quad \text{or} \quad u(t) = \sum_{\ell=0}^{+\infty} \frac{1}{\nu^\ell} U_\ell(\nu t)$$

where u_j , U_ℓ can be **computed recursively** (see in particular: Brachet et al., J. Fluid Mech., 1983 and J. Statist. Phys., 1984; Behr, Nečas and Wu, Math. Model. Numer. Anal., 2001; Sinai, J. Stat. Phys., 2005).

- Instead of discussing the difficult problem of convergence, one can consider the **truncated series**

$$u(t) = \sum_{j=0}^J u_j t^j \quad \text{or} \quad u(t) = \sum_{\ell=0}^L \frac{1}{\nu^\ell} U_\ell(\nu t)$$

and regard them as **approximate solutions**, to be analyzed with the method of **control inequalities**.

- For example, for the **BNW datum**, **extensive symbolic calculations** (by Python or Mathematica) allow to compute the previous truncated series for **Euler with $J \simeq 50$** , or **NS with $L \simeq 10$** (Morosi, Pernici and P, arXiv:1203.6865v1 [math.AP], 2012; Morosi and P, in preparation).

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- The analysis of the BNW datum via the above **truncated series** and **the control inequalities** **improves** the previous estimates based on the Galerkin approximation (existence now granted for Euler up to $t \simeq 0.2$; global existence granted for NS if $\nu \gtrsim 4$).

Here are some previous works related to this talk:

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