

Bounded vorticity, bounded velocity (Serfati) solutions to the incompressible 2D Euler equations

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Introduction

Equations for **incompressible, non-viscous (ideal) fluid flow in 2D** are given by:

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These are the **Euler equations** – they should be supplemented by initial data – $u(t=0) = u_0$ and – if there is a boundary – by slip boundary conditions $\rightarrow u \cdot \hat{n} = 0$ on the finite boundary, together with conditions at infinity.

Vorticity is defined by $\omega \equiv \partial_{x_1} u^2 - \partial_{x_2} u^1 \equiv \nabla^\perp \cdot u$, where $\nabla^\perp = (-\partial_{x_2}, \partial_{x_1})$. It is the only component of $\text{curl} (u^1, u^2, 0) \equiv (0, 0, \omega)$.

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Now, $K \in L_{loc}^p(\mathbb{R}^2)$, $1 \leq p < 2$ and K is q -th power integrable at infinity, with $q > 2$. Hence, to calculate $K * \omega$ need $\omega \in L^{p'} \cap L^{q'}$ with $p' > 2$ and $q' < 2$, e.g. $\omega \in L^\infty \cap L^1$.

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- $\omega_0 \in L_c^p$, $p > 1$, $u_0 \in L_{loc}^2$ then \exists (DiPerna and Majda 1987).
- $\omega_0 \in L_c^1$, $u_0 \in L_{loc}^2$ then \exists (Vecchi and Wu 1993).
- $\omega_0 \in \mathcal{BM}_{+,c} \cap H^{-1}$ then \exists (Delort 1991).

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This talk: discussion of Serfati's work, extension to continuous dependence on initial data and to flow domains exterior to a connected obstacle.

This is a report of **work in progress**.

Statement of problem

Let $D \subset \mathbb{R}^2$ be connected, open, bounded, smooth domain; $\Omega = \mathbb{R}^2 \setminus \overline{D}$.

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Fundamental questions: what is a ‘solution’? What happens to the Biot-Savart law? Where do we get ‘uniqueness’? What perturbations are allowed?

Related results

- Taniuchi (2004) gives complete, and very different proof of \exists , including Serfati flows, in full plane, allowing slightly unbounded initial vorticity (as Yudovich 1995 did in a bounded domain). Uses Littlewood-Paley decomposition and Bony's paradifferential calculus. Does not generalize to exterior domains.

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- Taniuchi, Tashiro, and Yoneda (2010) are concerned with *almost periodic flows in the full plane*. They prove \exists and ! assuming $u_0 \in L^\infty$ and $\omega_0 \in Y_{ul}^\theta$, for $\theta = \log(e + q)$; Y_{ul}^θ means "uniformly local" L^p -norms grow like $\theta(p)$ – includes Serfati initial data. Again, proof relies on Littlewood-Paley theory and Bony's paradifferential calculus; highly non-local proof.

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- Brunelli (2010) → studies full plane flows with velocity *growing* at infinity. He assumes $\omega_0 \in L^\infty$ and

$$\int \frac{1}{|x - y|} |\omega_0(y)| dy < \infty$$

for some $x \in \mathbb{R}^2$ and gets \exists and ! of (u, ω, Φ_t) such that $|u|$ grows at most like $\sqrt{|x|}$ at infinity. (Φ_t is the Lagrangian map.) The hypothesis excludes periodic flows.

Motivation

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- Local versus non-local;
- Need new idea to substitute Biot-Savart law;
- Broader potential applications in Serfati's key idea (new representation formula).

Serfati's representation formula

Key new idea: consider a smooth cutoff a_ε of the origin ($a_\varepsilon(z) = 1$ if $z \in \mathbb{R}^2$ and $|z|$ small, and vanishes if $|z|$ large) and establish a **formula** like

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$$u(t, x) = u_0(x) + \int a_\varepsilon(x - y)K(x - y)(\omega(t, y) - \omega_0(y)) dy + \\ - \int_0^t \int \left(u(s, y) \cdot \nabla_y \nabla_y [(1 - a_\varepsilon(x - y))K(x - y)]^\perp \right) \cdot u(s, y) dy ds.$$

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Indeed: $a_\varepsilon K \in L^1$ and $\omega \in L^\infty$ so the first integral converges; $\nabla_y \nabla_y [(1 - a_\varepsilon)K] \in L^1$ because of the extra decay at infinity coming from taking two derivatives, hence, if $u \in L^\infty$, then the second integral also converges.

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Hence,

$$\begin{aligned} \partial_t u &= \partial_t \int K(x - y)\omega(t, y) dy \\ &= \partial_t \int a_\varepsilon(x - y)K(x - y)\omega(t, y) dy + \int (1 - a_\varepsilon(x - y))K(x - y)\partial_t \omega dy. \end{aligned}$$

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Therefore, integrating in time yields

$$\begin{aligned} u(t, x) &= u_0(x) + \int a_\varepsilon(x - y)K(x - y)[\omega(t, y) - \omega_0(y)] dy \\ &\quad + \int_0^t \int (1 - a_\varepsilon(x - y))K(x - y)\partial_t \omega dy. \end{aligned}$$

Now,

$$\begin{aligned} \int (1 - a_\varepsilon(x - y))K(x - y)\partial_t\omega \, dy &= - \int (1 - a_\varepsilon(x - y))K(x - y)\mathbf{u} \cdot \nabla_y\omega \, dy \\ &= - \int (1 - a_\varepsilon(x - y))K(x - y) \operatorname{curl}_y(\mathbf{u} \cdot \nabla_y\mathbf{u}) \, dy. \end{aligned}$$

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Therefore, integrating by parts **carefully**, using $\operatorname{div} u = 0$, get

$$\begin{aligned} &\int (1 - a_\varepsilon(x - y))K(x - y)\partial_t\omega \, dy \\ &= \int \left(u(s, y) \cdot \nabla_y \nabla_y [(1 - a_\varepsilon(x - y))K(x - y)]^\perp \right) \cdot u(s, y) \, dy. \end{aligned}$$

Finally, substitute to obtain the desired **formula** – **Serfati identity**:

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$$\begin{aligned} &\int (1 - a_\varepsilon(x - y))K(x - y)\partial_t\omega \, dy \\ &= \int \left(u(s, y) \cdot \nabla_y \nabla_y [(1 - a_\varepsilon(x - y))K(x - y)]^\perp \right) \cdot u(s, y) \, dy. \end{aligned}$$

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$$\begin{aligned} u(t, x) &= u_0(x) + \int a_\varepsilon(x - y)K(x - y)[\omega(t, y) - \omega_0(y)] \, dy + \\ &- \int_0^t \int \left(u(s, y) \cdot \nabla_y \nabla_y [(1 - a_\varepsilon(x - y))K(x - y)]^\perp \right) \cdot u(s, y) \, dy. \end{aligned}$$

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5. Conclude u, ω satisfy incompressible 2D Euler in distributions, ω constant on particle paths and, also, representation formula remains valid. Also, (limit) u is log-Lipschitz.

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Hence, the *a priori* estimate becomes

$$\|u(t, \cdot)\|_{L^\infty} \leq \|u_0\|_{L^\infty} + C\varepsilon\|\omega_0\|_{L^\infty} + \frac{C}{\varepsilon} \int_0^t \|u(s, \cdot)\|_{L^\infty}^2 ds.$$

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The desired L^∞ -estimate for u follows, hence, by Gronwall's lemma.

Exterior domain

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Second, (hindsight from previous exterior domain work) know

$$K_{\Omega}(x, y) + K(x) \sim K(x - y)$$

($K(x) = x^{\perp}/(2\pi|x|^2)$), not $K_{\Omega} \sim K(x - y)$.

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 u(t, x) &= u_0(x) + \int_{\Omega} a_{\varepsilon}(x - y) J(x, y) (\omega(t, y) - \omega_0(y)) dy \\
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 &\quad + \frac{K(x)}{2} \int_0^t \int_{\Gamma} |u(s, y)|^2 \nabla a_{\varepsilon}(x - y) \cdot d\sigma(y) ds.
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Uniqueness

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Introduce the Serfati space \mathcal{S} , of divergence-free vector fields tangent to the boundary with the norm $\|u\|_{\mathcal{S}} = \|u\|_{L^\infty} + \|\operatorname{curl} u\|_{L^\infty}$.

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Theorem

Let Ω be a smooth domain exterior to a connected and bounded set. Let $u_0 \in S$. Then there exists one and at most one weak solution of Euler in Ω with initial velocity u_0 .

Continuous dependence

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Theorem

Suppose initial velocities, u_1^0 and u_2^0 , with vorticities, ω_1^0 and ω_2^0 , such that $u_1^0 - u_2^0$ lies in

$$S_p := \{u \in (L^\infty(\Omega))^2 : \operatorname{div} u = 0, u \cdot \mathbf{n} = 0, \omega \in L^p(\Omega)\}$$

for some p in $(2, \infty]$, with $\|\cdot\|_{S^p} = \|\cdot\|_{L^\infty} + \|\omega(\cdot)\|_{L^p}$. Then, for all sufficiently small $s_0 = \|u_1^0 - u_2^0\|_{S^p}$,

$$\|u_1(t) - u_2(t)\|_{L^\infty} \leq s_0 e^{Ct} + C_t (s_0 t) e^{-C_t(2+t)} [\log C_t + s_0 t e^{-C_t(2+t)}] \\ [C(2+t)e^{Ct} + 1],$$

where C and C_t depend on the initial data and on p , with C_t a continuous function of time.

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2. Should try to improve continuous dependence result so as to have same norm comparison, not velocity stable in L^∞ in terms of initial perturbation in S . Recall Taniuchi et alli, continuous dependence in $B_{\infty,1}^0$. But only for full plane and only by Littlewood-Paley...

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 - (ii) $u(x) = x^\perp/|x|$ has vorticity $1/|x|$, hence Serfati. Note that Biot-Savart law cannot be used.
 - (iii) Serfati vorticity with no decay: vorticity is characteristic function of a strip lying outside/”above” obstacle, velocity interpolates linearly from 0 ”below” strip to constant parallel to and ”above” strip. For example, if the strip is $\{2 < x_2 < 3\}$ then

3. One natural question: can we characterize those *vorticities* which “are Serfati”? I.e., bounded vorticities which are the curl of a bounded velocity? Note that, if $\omega_0 \equiv 1$ then u_0 is not bounded, hence it is *not* Serfati.
4. Some examples of Serfati flows:
- (i) All smooth (doubly) periodic flows are Serfati. Adding a compactly supported function to the vorticity does not change this.
 - (ii) $u(x) = x^\perp/|x|$ has vorticity $1/|x|$, hence Serfati. Note that Biot-Savart law cannot be used.
 - (iii) Serfati vorticity with no decay: vorticity is characteristic function of a strip lying outside/”above” obstacle, velocity interpolates linearly from 0 ”below” strip to constant parallel to and ”above” strip. For example, if the strip is $\{2 < x_2 < 3\}$ then

$$u = u(x) = \begin{cases} (1, 0) & \text{if } x_2 > 3, \\ (x_2 - 2, 0) & \text{if } 2 < x_2 < 3, \\ (0, 0) & \text{if } x_2 < 2, \end{cases}$$

is Serfati in the exterior of the unit disk.

Thank you!