

# Blow up at the hyperbolic boundary for a $2 \times 2$ system arising from chemical engineering

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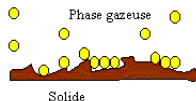
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26 Juin 2012

- 1 Introduction
  - Chromatography
  - The model
  - Sorption effect
- 2 Gas-chromatography with sorption effect for two species
  - Hyperbolicity, Riemann invariants
  - Entropies
  - Riemann Problem
- 3 Existence of entropy solutions
  - Godunov scheme
  - Front Tracking Algorithm (FTA)
- 4 An example of “Blow up”
- 5 Prospects and open problems
- 6 Bibliography

- 1 Introduction
  - Chromatography
  - The model
  - Sorption effect
- 2 Gas-chromatography with sorption effect for two species
  - Hyperbolicity, Riemann invariants
  - Entropies
  - Riemann Problem
- 3 Existence of entropy solutions
  - Godunov scheme
  - Front Tracking Algorithm (FTA)
- 4 An example of “Blow up”
- 5 Prospects and open problems
- 6 Bibliography

## Chromatography gas-solid : a technic for analyzing a gaseous mixture

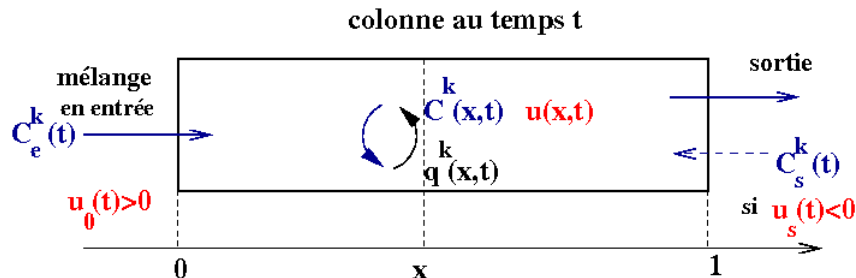


The mixture analysed is vaporised at the entrance of a column that contains a solid substance (the absorber) called the stationary phase and then it is transported across it by a gas carrier. The gas carrier (or gas vector) is the mobile phase. In most cases it has to be inerted vis a vis the solutes and the stationary phase.

# Unknowns

During the process,  $d$  species exist together under 2 phases

- one gas phase of concentration  $c_i$  and moving with speed  $u$
- one solid (adsorbed) with concentration  $q_i^*$



We define by  $u$  the speed of a tracer that is inert in the column. This is a measure of the speed of the flow across the column. **In general this speed can not be considered constant.**

## Chemical assumptions

- constant temperature
- 1D model,  $x \in \mathbb{R}_+$
- negligible axial diffusion
- perfect gas
- instantaneous exchange

## Prospects

- nonisothermal model
- validity of the model
- relaxation

Mass conservation :

$$\partial_t(c_i + \mathbf{q}_i^*) + \partial_x(u c_i) = 0, \quad 1 \leq i \leq d$$

$\mathbf{q}_i^*$  depends on all  $c_j$  :

$$\mathbf{q}_i^* = \mathbf{q}_i^*(c_1, \dots, c_d)$$

$\mathbf{q}_i^*$  is the  $i^{\text{th}}$  isothermal law given by the experiments (chemistry)

# Classic isotherms

**Linear isotherm** :  $q_i^* = K_i c_i$ , avec  $K_i \geq 0$

**Langmuir isotherm** :

$$q_i^* = \frac{Q_i K_i c_i}{1 + \sum_{j=1}^d K_j c_j}, \text{ with } K_i \geq 0, Q_i > 0$$

**BET isotherm** : one active gas and one inert gas ( $q_2^* = 0$ )

$$q_1^* = \frac{Q K c_1}{(1 + K c_1 - (c_1/c_s))(1 - (c_1/c_s))}, \quad Q > 0, K > 0, c_s > 0,$$



Unknowns : velocity  $u$ , gaseous concentrations  $c_j$

**In gas-solid chromatography, the velocity  $u$  is not constant** : it depends on the mixture composition, which depends on mass transfer between phases.

The velocity of the mixture has to be found in order to achieve a given pressure (or density in the isothermal model). The sorption effect is taken into account through a constraint

$$\sum_{i=1}^d c_i = \rho(t)$$

$\rho$  : the given total density of the mixture (the same in the column)

- 1 Introduction
  - Chromatography
  - The model
  - Sorption effect
- 2 Gas-chromatography with sorption effect for two species
  - Hyperbolicity, Riemann invariants
  - Entropies
  - Riemann Problem
- 3 Existence of entropy solutions
  - Godunov scheme
  - Front Tracking Algorithm (FTA)
- 4 An example of “Blow up”
- 5 Prospects and open problems
- 6 Bibliography

$\rho \equiv 1$  and  $x \in \mathbb{R}_+$  :

$$\partial_t(c_1 + \mathbf{q}_1^*(c_1, c_2)) + \partial_x(u c_1) = 0 \quad (1)$$

$$\partial_t(c_2 + \mathbf{q}_2^*(c_1, c_2)) + \partial_x(u c_2) = 0 \quad (2)$$

$$c_1 + c_2 = 1 \quad (3)$$

Let  $c = c_1$  and  $\mathbf{q}_i(c) = \mathbf{q}_i^*(c, 1 - c)$   $i = 1, 2$

(1) + (2) yields with (3) :

$$\partial_t(\mathbf{q}_1(c) + \mathbf{q}_2(c)) + \partial_x u = 0$$

**N.B.** : isotherm properties :  $q'_1 = \frac{dq_1}{dc} \geq 0$  and  $q'_2 = \frac{dq_2}{dc} \leq 0$

# The $2 \times 2$ system

$$\left\{ \begin{array}{l} \partial_t(c + q_1(c)) + \partial_x(uc) = 0 \\ \partial_t(\underbrace{q_1(c) + q_2(c)}_{= h(c)}) + \partial_x u = 0 \end{array} \right.$$

**Initial boundary values :**

$$\left\{ \begin{array}{l} c(0, x) = c_0(x) \in [0, 1], \quad x > 0 \\ c(t, 0) = c_b(t) \in [0, 1], \quad t > 0 \\ u(t, 0) = u_b(t) > 0, \quad t > 0 \end{array} \right.$$

# The $2 \times 2$ system

We analyse the system in terms of hyperbolicity system of P.D.E provided we exchange the time and space variables :

**$x$  is the evolutive variable, not  $t$**

$$\begin{cases} \partial_x(u c) + \partial_t(c + q_1(c)) = 0 \\ \partial_x u + \partial_t h(c) = 0 \end{cases}$$

**with conservative quantities**

$$m = u c, \quad u = u \rho$$

$m$  is the flow rate of the first species

$u\rho$  is the total flow rate

$$\text{Eigenvalues : } \mathbf{0} \text{ and } \lambda = \frac{H(c)}{u}$$

$$\text{with } H(c) = 1 + (1 - c) q'_1 - c q'_2 \geq 1$$

If  $u > 0$ , the system is strictly hyperbolic. Moreover :

$$d\lambda \cdot r = \frac{H(c)}{u^2} f''(c)$$

with  $r$  the associated eigenvector to  $\lambda$  and

$$\mathbf{f}(\mathbf{c}) = q_1(c) - c h(c) = (1 - c)q_1(c) - cq_2(c)$$

**$\lambda$  is genuinely nonlinear if  $f'' \neq 0$**

## Langmuir isotherm :

$$q_1(c) = \frac{Q_1 K_1 c}{1 + K_1 c + K_2 (1 - c)}, \quad q_2(c) = \frac{Q_2 K_2 (1 - c)}{1 + K_1 c + K_2 (1 - c)}$$

**$f''$  keeps a constant sign**

**BET isotherm :** with an inert gas ( $q_2(c) = 0$ ) and

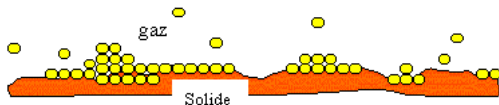
$$q_1(c) = \frac{Q K c}{(1 + Kc - (c/c_s))(1 - (c/c_s))}, \quad Q > 0, K > 0, c_s > 0,$$

**then  $f(c) = (1 - c) q_1(c)$  is not concave and not convex**

Generally, isotherms are not convex except some important cases  
(Langmuir, ammoniac, water vapour)

## Interpretation :

at each inflexion point a new layer starts on the pervious layer





For smooth solutions

$$\partial_x (u G(c)) = 0$$

$$\partial_x c + \frac{H(c)}{u} \partial_t c = 0$$

with

$$g' = \frac{-h'}{H}, \quad G = \exp(g)$$

There are two Riemann invariants :

$$c \quad \text{and} \quad W = u G(c) = u e^{g(c)}$$

Entropies :

$$\mathbf{S}(\mathbf{c}, u) = \phi(u\mathbf{G}(\mathbf{c})) + u\psi(\mathbf{c})$$

where  $\phi$  and  $\psi$  are smooth functions

Entropy-flux  $Q = Q(\mathbf{c})$  satisfies

$$\mathbf{Q}'(\mathbf{c}) = \mathbf{h}'(\mathbf{c})\psi(\mathbf{c}) + H(\mathbf{c})\psi'(\mathbf{c})$$

$$\Rightarrow \partial_t \mathbf{S} + \partial_x Q = 0$$

# Convex entropy ?

For any **convex** function  $\psi$ ,

$$\mathbf{S}_2(\mathbf{c}, u) = u \psi(\mathbf{c})$$

is convex (but not strictly convex)

There are strictly convex entropies of the form

$$\mathbf{S}_1(\mathbf{c}, u) = \phi(uG(\mathbf{c})) \Leftrightarrow$$

$$G'' = (g'' + g'^2) \exp(g) \neq 0 \text{ for } c \in [0, 1]$$

## Definition

Let  $T > 0$ ,  $X > 0$ ,

$u \in L^\infty((0, T) \times (0, X))$  and  $0 \leq c(t, x) \leq 1$  p.p. in  $(0, T) \times (0, X)$

then  $(c, u)$  is a weak entropy solution if for any convex function  $\psi$  :

$$\partial_x (u \psi(c)) + \partial_t Q(c) \leq 0$$

with  $Q' = H\psi' + h'\psi$

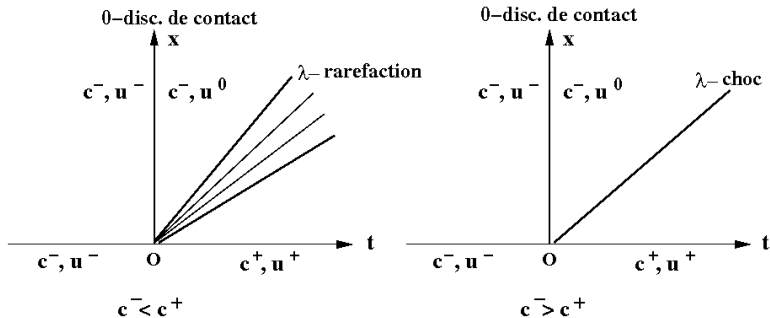
**Remark :** if  $\pm G'' > 0$  on  $[0, 1]$ , we also have :

$$\pm \partial_x (u G(c)) \leq 0, \text{ on } [0, 1]$$

# Riemann problem

Rappel :  $c$  is a Riemann invariant

If  $f$  is convex :



through a  $\lambda$ -wave  $TV[\ln(u(z))] \leq \gamma |c^+ - c^0|, z = \frac{x}{t}$

- 1 Introduction
  - Chromatography
  - The model
  - Sorption effect
- 2 Gas-chromatography with sorption effect for two species
  - Hyperbolicity, Riemann invariants
  - Entropies
  - Riemann Problem
- 3 Existence of entropy solutions
  - Godunov scheme
  - Front Tracking Algorithm (FTA)
- 4 An example of “Blow up”
- 5 Prospects and open problems
- 6 Bibliography

The convergence of the Godunov scheme is generally an open problem for a  $2 \times 2$  system

- LeVeque, Randall J. ; Temple, Blake, 1985  
Stability of Godunov's method for a class of  $2 \times 2$  systems of conservation laws  
Trans. Amer. Math. Soc. 288
- Bressan, Alberto ; Jenssen, Helge Kristian, 2001  
Convergence of the Godunov scheme for straight line systems
- Bressan, Alberto ; Jenssen, Helge Kristian ; Baiti, Paolo, 2006  
An instability of the Godunov scheme  
Comm. Pure Appl. Math. 59

## An existence result

Let be  $X > 0$ ,  $T > 0$ . We assume :

$$c_0 \in BV(0, X), c_b \in BV(0, T), u_b \in L^\infty(0, T)$$

$$0 \leq c_0, c_b \leq 1 \quad \text{and} \quad \inf_{0 < t < T} u_b(t) > 0$$

Then there exists an entropy solution  $(c, u)$  such that :

$$0 \leq \min(\inf c_b, \inf c_0) \leq c \leq \max(\sup c_b, \sup c_0) \leq 1$$

$$|u| \leq \|u_b\|_\infty \exp(\gamma TV(c)) \quad \text{and} \quad \inf u > 0$$

$$c \in L^\infty((0, T); BV(0, X)) \cap L^\infty((0, X); BV(0, T))$$

$$\ln(u) \in L^\infty((0, T); BV(0, X))$$



# Front Tracking Algorithm (FTA)

With the assumption :

the eigenvalue  $\lambda$  is genuinely non linear

this scheme shows again the existence of an entropy solution but more precisely the specific structure of the velocity in two cases :

- the case with smooth concentrations
- the more complicated case with  $BV$  concentrations

$$u_b \in BV[0, T] \Rightarrow u \in BV([0, T] \times [0, X])$$

not given by the Godunov scheme !

A. Corli and O. Gues, Stratified solutions for systems of conservation laws, Trans. Amer. Math. Soc., 2001

# Front Tracking Algorithm (FTA)

The stratified specific structure of the velocity  $L_t^\infty \times BV_{t,x}$

$$u(t, x) = u_b(t)v(t, x)$$

has interesting applications

- the stability in strong topologies of concentrations with respect to weak\* limits for the incoming velocity
- an example of **blow-up**

but there is an assumption  $f'' > 0$  and then the eigenvalue  $\lambda$  is genuinely non linear

# Contents

- 1 Introduction
  - Chromatography
  - The model
  - Sorption effect
- 2 Gas-chromatography with sorption effect for two species
  - Hyperbolicity, Riemann invariants
  - Entropies
  - Riemann Problem
- 3 Existence of entropy solutions
  - Godunov scheme
  - Front Tracking Algorithm (FTA)
- 4 An example of “Blow up”
- 5 Prospects and open problems
- 6 Bibliography

## Examples for $N \times N$ hyperbolic systems

- with  $N \geq 3$ 
  - Helge Kristian **Jenssen** ; Carlo Sinestrari. CPDE (1999)
  - Helge Kristian **Jenssen** ; Robin **Young**. JHPDE (2004)
  - Michael, Sever  
Distribution solutions of nonlinear systems of conservation laws  
Mem. Amer. Math. Soc. 190 (2007),
- with  $N = 2$ 
  - Robin **Young** : SIMA 99, CM 03, CMS 03  
**2 Burgers equations linearly coupled at 2 boundaries**
  - **Bourdarias, Gisclon, Junca, JHPDE (2010).**  
Blow up for a  $2 \times 2$  chemical system with **1 characteristic boundary**

- Maximum Principle for  $c(t, x) : 0 \leq c \leq 1$  but not for  $u(t, x)$ !

$$u \rightarrow +\infty$$

- $G'' < 0 \Rightarrow \partial_x(uG(c)) \geq 0$
- avoid Temple system
  - Temple, Blake, Trans. Amer. Math. Soc. (1983)
  - Bressan, Alberto ; Goatin, Paola  
Stability of  $L^\infty$  solutions of Temple class systems  
Differential Integral Equations 13 (2000)
  - Bianchini, Stefano  
Stability of  $L^\infty$  solutions for hyperbolic systems with coinciding shocks and rarefactions  
SIAM J. Math. Anal. 33 (2001)

# Temple system ?

Generally, it is not a Temple system (ammonia, water vapor)

If  $f'' > 0$ , it is a Temple system iff

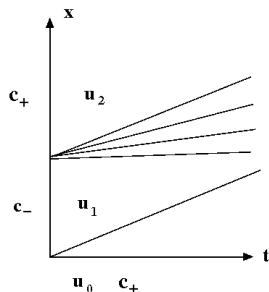
$$\partial_x(uG(c)) = 0$$

## Examples :

If  $G'' \equiv 0$  (for example two linear isotherms) then it is a Temple system

If  $q_1 \equiv 0$  (gas 1 inert), it is a Temple system iff  $q_2'' \equiv 0$

# Brick for the blow up



We assume that the system is not a Temple system, that

$$f'' < 0 \text{ and } h' < 0 \text{ (no restrictive)}$$

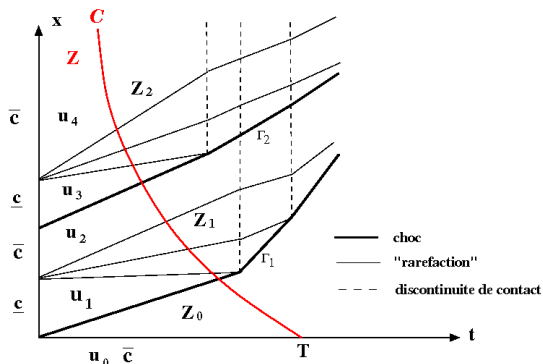
then there exists a choice of  $c_- < c_+$  such that

$$u_2 = \mathbf{R}(c_-, c_+) u_0$$

with

$$G'' < 0 \Rightarrow W = uG(c) \nearrow \Rightarrow \mathbf{R}(c_-, c_+) \geq 1$$

## Iterations



- $\forall (c_-, c_+) \quad R(c_-, c_+) = 1 \Leftrightarrow$  the system is in the Temple class
- if the system is not a Temple system then

$$\exists c_- < c_+, \quad \mathbf{R}(c_-, c_+) > 1$$



Blow up at  $X > 0$

- In  $Z$ , explicit solution before interactions
- Interactions outside  $Z$
- **Unique** piecewise smooth solution on  $[0, X - \varepsilon] \times \mathbb{R}$ ,  $\varepsilon > 0$   
Bourdarias, Gisclon, Junca : JMAA 06
- $u(0, X) = +\infty$
- The boundary becomes twice characteristic :  $\lambda = 0$
- $TVc_0 = TVc(t = 0, [0, X]) = +\infty$

# An example of “Blow up”

## Blow-up

Under the assumptions :

- $G'' < 0$
- $h'$  et  $f''$  do not vanish
- the system is not a Temple system

we can construct an solution such that  $\forall T > 0, \forall X > 0$ , when  $x \rightarrow X$

$\|u(\cdot, x)\|_{L^\infty(0,T)} \rightarrow +\infty, \lambda = \frac{H(c)}{u} \rightarrow 0, c$  stays bounded

Remarks :

- $G'' < 0 \Rightarrow \partial_x(uG(c)) \geq 0$  :  $-uG(c)$  is a convex entropy
- $h' \neq 0 \Rightarrow$  one gas is more active than an another
- $f'' \neq 0 \Rightarrow$  the eigenvalue  $\lambda$  is genuinely non linear

Chemical Examples : [ammoniac](#), [water vapor](#)

# Contents

- 1 Introduction
  - Chromatography
  - The model
  - Sorption effect
- 2 Gas-chromatography with sorption effect for two species
  - Hyperbolicity, Riemann invariants
  - Entropies
  - Riemann Problem
- 3 Existence of entropy solutions
  - Godunov scheme
  - Front Tracking Algorithm (FTA)
- 4 An example of “Blow up”
- 5 Prospects and open problems
- 6 Bibliography

# Prospects and open problems

- Existence of solution on  $[0, X = +\infty[$  with

$$u(0, x) = +\infty \text{ for } x \geq X?$$

(proof tentative : problem with the O-waves)

- Validity of the model
- Relaxation
- Theory  $L^\infty$  with some restrictions on the data  
admissible initial boundary value data  $u_b, c_b, c_0 \in L^\infty$

$$\text{such that } u(0, X) < +\infty$$

- Existence in fractional  $BV$  spaces :  $BV^s$ ,  $0 < s < 1$  with

$$u_b \in L^\infty \text{ and } c_0, c_b \in BV^{1/2} \text{ or } 1/3$$

- Kinetic problem
- Nonisothermal model

# Contents

- 1 Introduction
  - Chromatography
  - The model
  - Sorption effect
- 2 Gas-chromatography with sorption effect for two species
  - Hyperbolicity, Riemann invariants
  - Entropies
  - Riemann Problem
- 3 Existence of entropy solutions
  - Godunov scheme
  - Front Tracking Algorithm (FTA)
- 4 An example of “Blow up”
- 5 Prospects and open problems
- 6 Bibliography

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- *Existence of Weak Entropy Solutions for Gas Chromatography system with one or two active species and non Convex Isotherms.* Commun. Math. Sci., 2007
- *Hyperbolic models in gas-solid chromatography* Bol. Esp. Mat. Apl., 2008
- *Strong Stability with respect to weak limit for a Hyperbolic System arising from Gas Chromatography.* Methods Appl. Anal., 2010
- *Blow up at the hyperbolic boundary for a  $2 \times 2$  system arising from chemical engineering.* J. Hyperbolic Differ. Equ., 2010
- *$BV^s$  spaces and applications to scalar conservation laws*
- *A kinetic scheme for a hyperbolic system arising in gas chromatography*